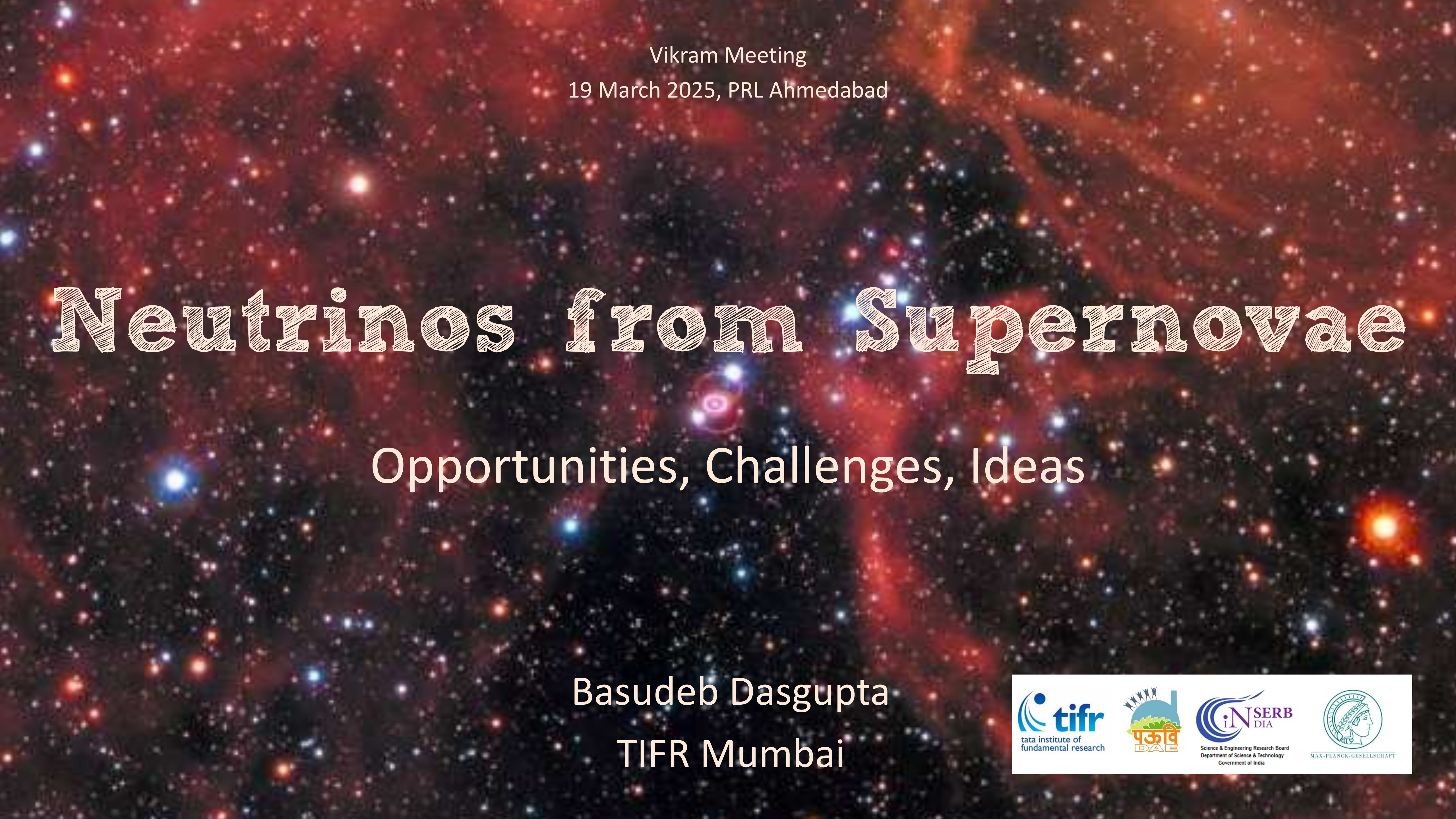


Vikram Meeting

19 March 2025, PRL Ahmedabad

Neutrinos from Supernovae

Opportunities, Challenges, Ideas

Basudeb Dasgupta

TIFR Mumbai

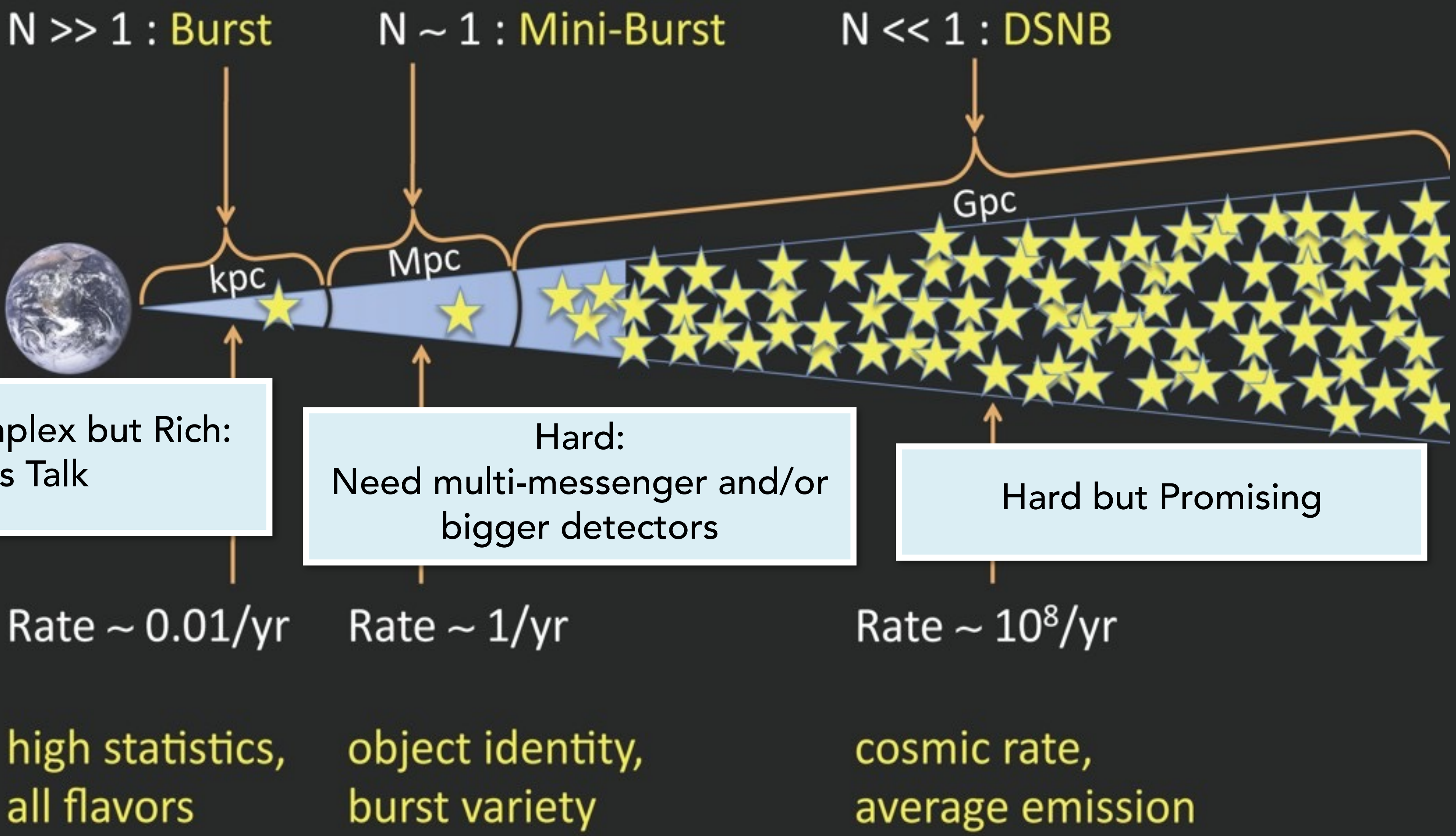
tifr
tata institute of fundamental research

पऊवि
DAE

SERB
INDIA
Science & Engineering Research Board
Department of Science & Technology
Government of India

MAX-PLANCK-GESELLSCHAFT

Here, There, Everywhere



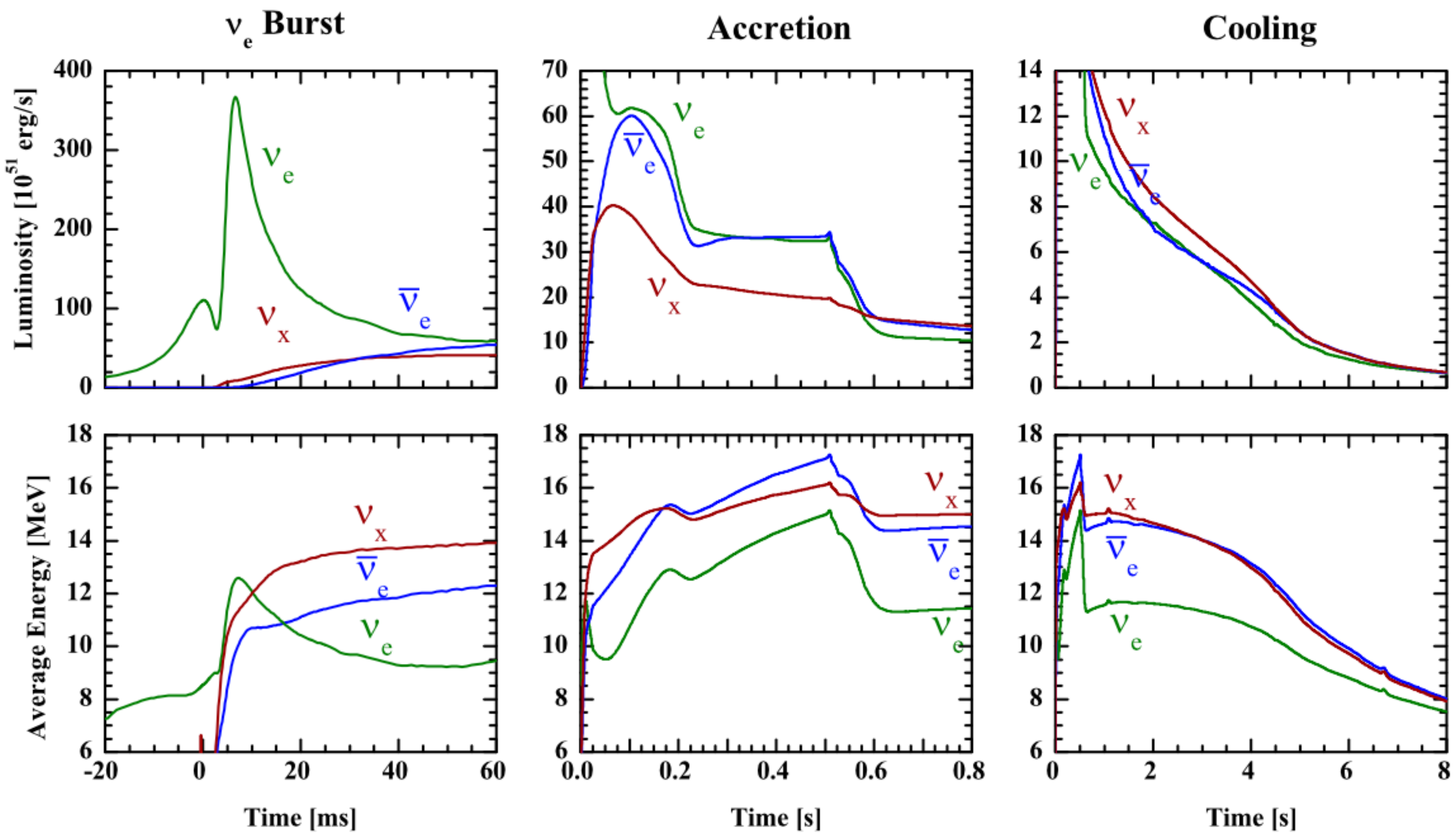
Rare & Complex but Rich:
This Talk

Hard:
Need multi-messenger and/or
bigger detectors

Hard but Promising

John Beacom, Ohio State University Neutrino 2012, Kyoto, Japan, June 2012

Neutrinos from a SN



1d simulation of a 27 M_{sun} star by Garching group
See review by Janka, Melson, and Summa (2016)

Modelling is driven by SN simulations, though most features are theoretically well- justified

SN1987A is the only data.

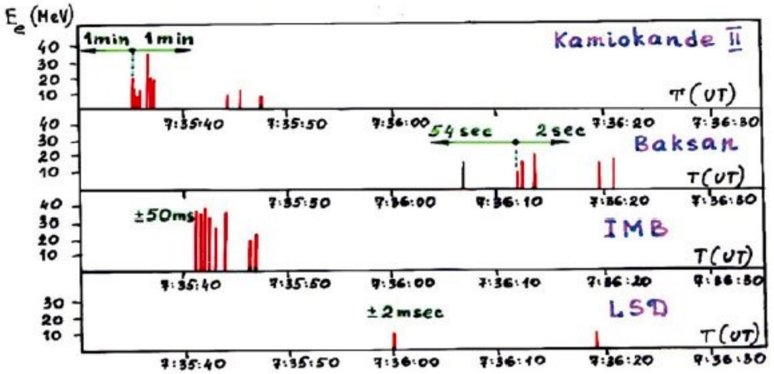
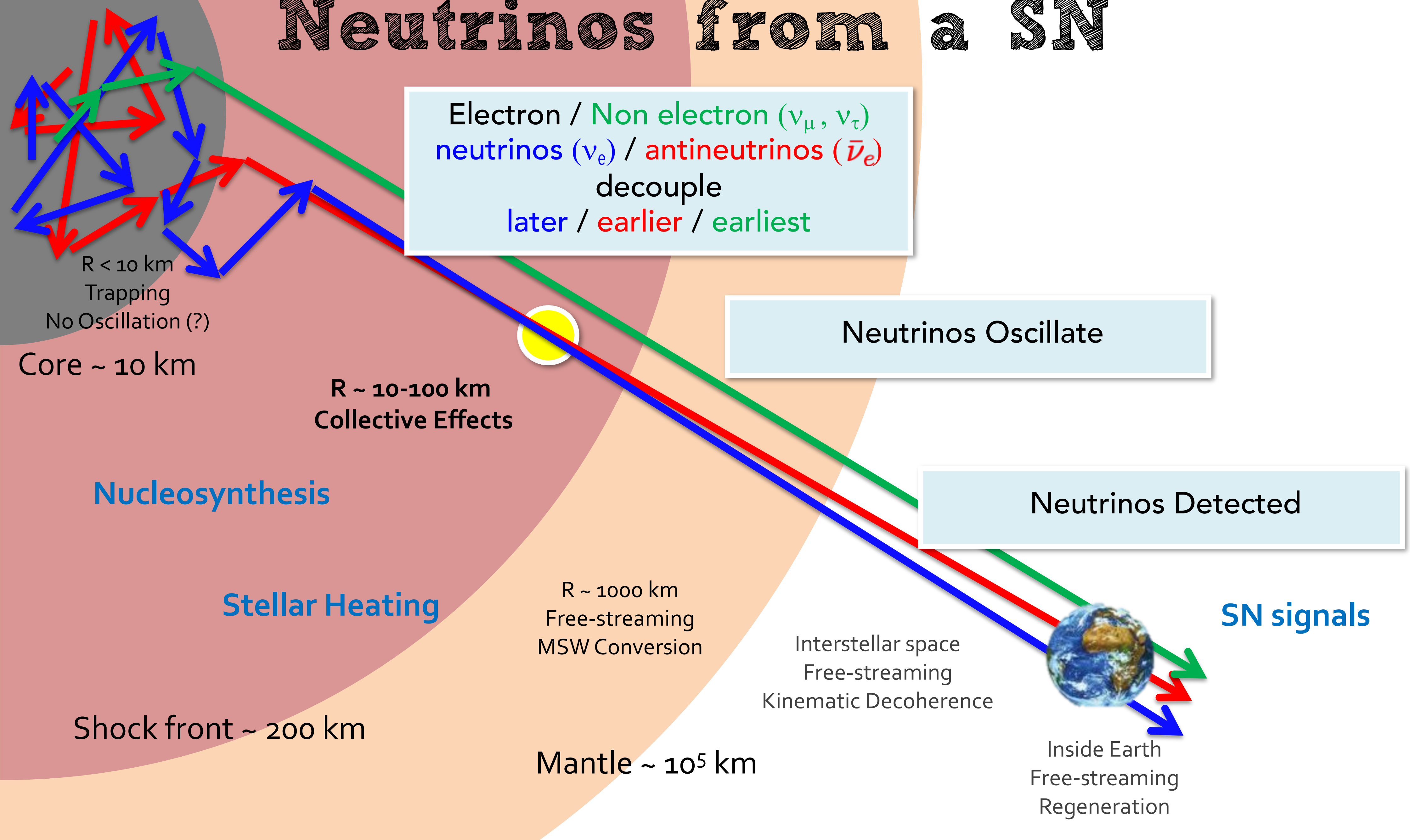


Image credit: Suzuki (2003)

Neutrinos from a SN



Electron / Non electron (ν_μ , ν_τ)
neutrinos (ν_e) / antineutrinos ($\bar{\nu}_e$)
decouple
later / earlier / earliest

Neutrinos Oscillate

Neutrinos Detected

R < 10 km
Trapping
No Oscillation (?)

Core ~ 10 km

R ~ 10-100 km
Collective Effects

Nucleosynthesis

Stellar Heating

Shock front ~ 200 km

R ~ 1000 km
Free-streaming
MSW Conversion

Mantle ~ 10^5 km

Interstellar space
Free-streaming
Kinematic Decoherence

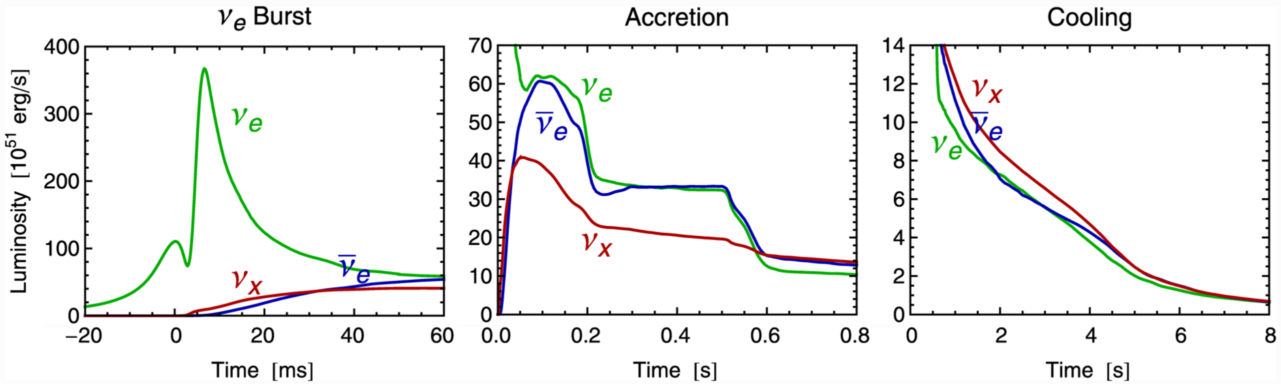
SN signals

Inside Earth
Free-streaming
Regeneration

SN Neutrino Program

	Detector	Type	Mass (kt)	Location	Events	Flavors
Running	Super-Kamiokande	H_2O	32	Japan	7,000	$\bar{\nu}_e$
	LVD	C_nH_{2n}	1	Italy	300	$\bar{\nu}_e$
	KamLAND	C_nH_{2n}	1	Japan	300	$\bar{\nu}_e$
	Borexino	C_nH_{2n}	0.3	Italy	100	$\bar{\nu}_e$
	IceCube	Long string	(600)	South Pole	(10^6)	$\bar{\nu}_e$
	Baksan	C_nH_{2n}	0.33	Russia	50	$\bar{\nu}_e$
	MiniBooNE*	C_nH_{2n}	0.7	USA	200	$\bar{\nu}_e$
	HALO	Pb	0.08	Canada	30	ν_e, ν_x
	Daya Bay	C_nH_{2n}	0.33	China	100	$\bar{\nu}_e$
	NOνA*	C_nH_{2n}	15	USA	4,000	$\bar{\nu}_e$
Future	SNO+	C_nH_{2n}	0.8	Canada	300	$\bar{\nu}_e$
	MicroBooNE*	Ar	0.17	USA	17	ν_e
	DUNE	Ar	34	USA	3,000	ν_e
	Hyper-Kamiokande	H_2O	560	Japan	110,000	$\bar{\nu}_e$
	JUNO	C_nH_{2n}	20	China	6000	$\bar{\nu}_e$
	RENO-50	C_nH_{2n}	18	Korea	5400	$\bar{\nu}_e$
	LENA	C_nH_{2n}	50	Europe	15,000	$\bar{\nu}_e$
	PINGU	Long string	(600)	South Pole	(10^6)	$\bar{\nu}_e$

From review by Scholberg (2012)



Burst	Accretion	Cooling
SN standard candle	Collective effects	Nuclear physics
Mass ordering	SN theory	Exotics/Axions
Timing	Mass ordering	Nucleosynthesis
…	Pointing	Shock
	…	…

Several detectors capable of detecting SN neutrinos.
There is a rich science-case.

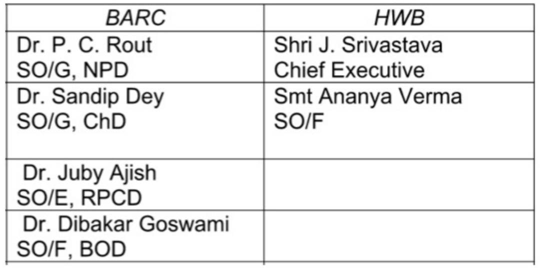
Deuterated Liquid Scintilator

We propose a kton-scale deuterated liquid scintillator with added Gd, and instrumented with PMTs that can be used to study low energy neutrinos esp. through the Neutral Current channel.

The Road to a Deuterium-based Scintillator Detector

based on the DLS Science Case White Paper
+ JCAP 11 (2021) 005 by Bhavesh Chauhan, BD, Vivek Datar
+ ongoing work by the DLS Study Group at TIFR
+ ongoing work by DLS Task Force of DAE

DAE Task Force, TIFR Study Group



BARC	*HWB*
Dr. P. C. Rout SO/G, NPD	Shri J. Srivastava Chief Executive
Dr. Sandip Dey SO/G, ChD	Smt Ananya Verma SO/F
Dr. Juby Ajish SO/E, RPCD	
Dr. Dibakar Goswami SO/F, BOD	

Prof. Vivek M. Datar, IMSc
Prof. D. Indumathi, IMSc
Prof. Milind Diwan, Brookhaven National Lab, USA
Prof. Deepak Samuel, Central University of Karnataka
Prof. Basudeb Dasgupta, TIFR
Prof. Amol S. Dighe, TIFR
Prof. Gobinda Majumder, TIFR
Prof. Kajari Mazumdar, TIFR
Prof. Vandana S. Nanal, TIFR
Prof. Rudrajyoti Palit, TIFR

... and several others ...
(e.g., participants of the Underground Science Lab meeting held in 2022 at TIFR)

Deuterated Detectors

Charged Current

Flavor	Ordinary Detector		Deuterated Detector	
	Channel	Detector	Channel	Detector
ν_e	ν_e Ar $\rightarrow e^-$K*	DUNE *Spectrum* ✓ *Tagging* ✗	$\nu_e\ d \rightarrow e^-\ p\ p$	SNO *Spectrum* ✓ *Tagging* ✗
	ν_e Pb $\rightarrow \nu_e$ Bi*	HALO *Spectrum* ✗ *Tagging* ✗	$\nu_e\ d \rightarrow e^-\ p\ p$	DLS *Spectrum* ✓ *Tagging* ✓
$\bar{\nu}_e$	$\bar{\nu}_e\ p \rightarrow e^+\ n$	SuperK+Gd *Spectrum* ✓ *Tagging* ✓	$\bar{\nu}_e d \rightarrow e^+\ n\ n$	SNO *Spectrum* ✓ *Tagging* ✓
	$\bar{\nu}_e\ p \rightarrow e^+\ n$	LVD, JUNO *Spectrum* ✓ *Tagging* ✗	$\bar{\nu}_e d \rightarrow e^+\ n\ n$	DLS *Spectrum* ✓ *Tagging* ✓

Neutral Current

Flavor	Ordinary Detector		Deuterated Detector	
	Channel	Detector	Channel	Detector
$\overset{(-)}{\nu}$	$\overset{(-)}{\nu} p \rightarrow \overset{(-)}{\nu} p$	JUNO, THEIA *Spectrum* ✓ *Tagging* ✗	$\overset{(-)}{\nu} d \rightarrow \overset{(-)}{\nu}\ p\ n$	SNO *Spectrum* ✗ *Tagging* ✗
	$\overset{(-)}{\nu}$ Pb $\rightarrow \overset{(-)}{\nu}$ Pb*	HALO *Spectrum* ✗ *Tagging* ✗	$\overset{(-)}{\nu} d \rightarrow \overset{(-)}{\nu}\ p\ n$	DLS *Spectrum* ✓ *Tagging* ✓

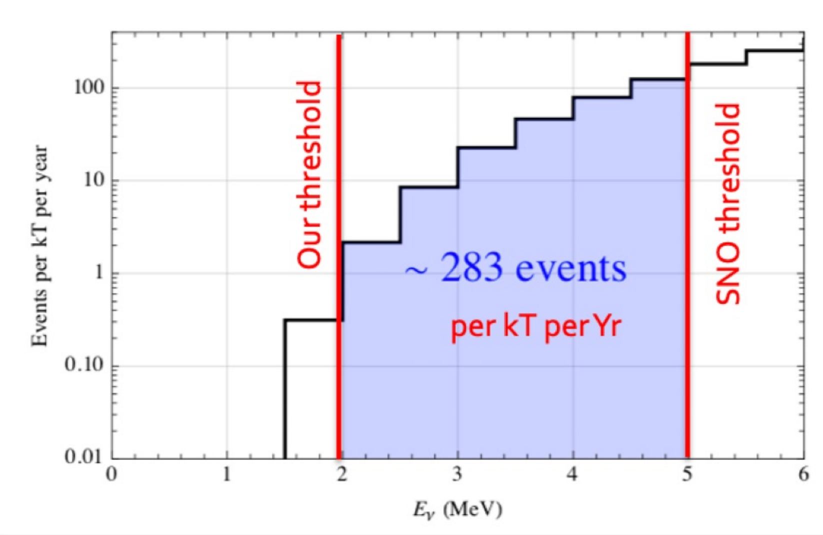
Spectrum ~ "Can reconstruct incident neutrino spectrum"

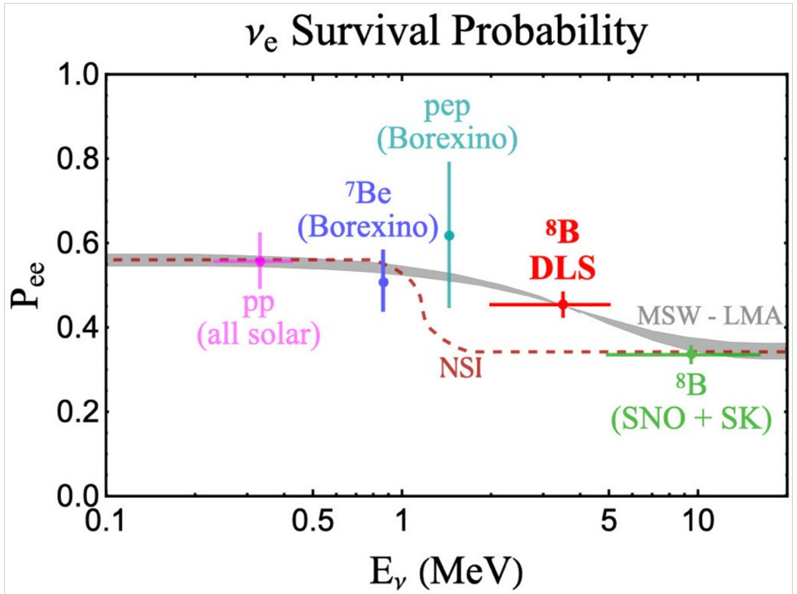
Tagging ~ "Multiple detectable particles in the final state"

Significant Advantages of Multiple Detection Channels

Deuterated Detectors

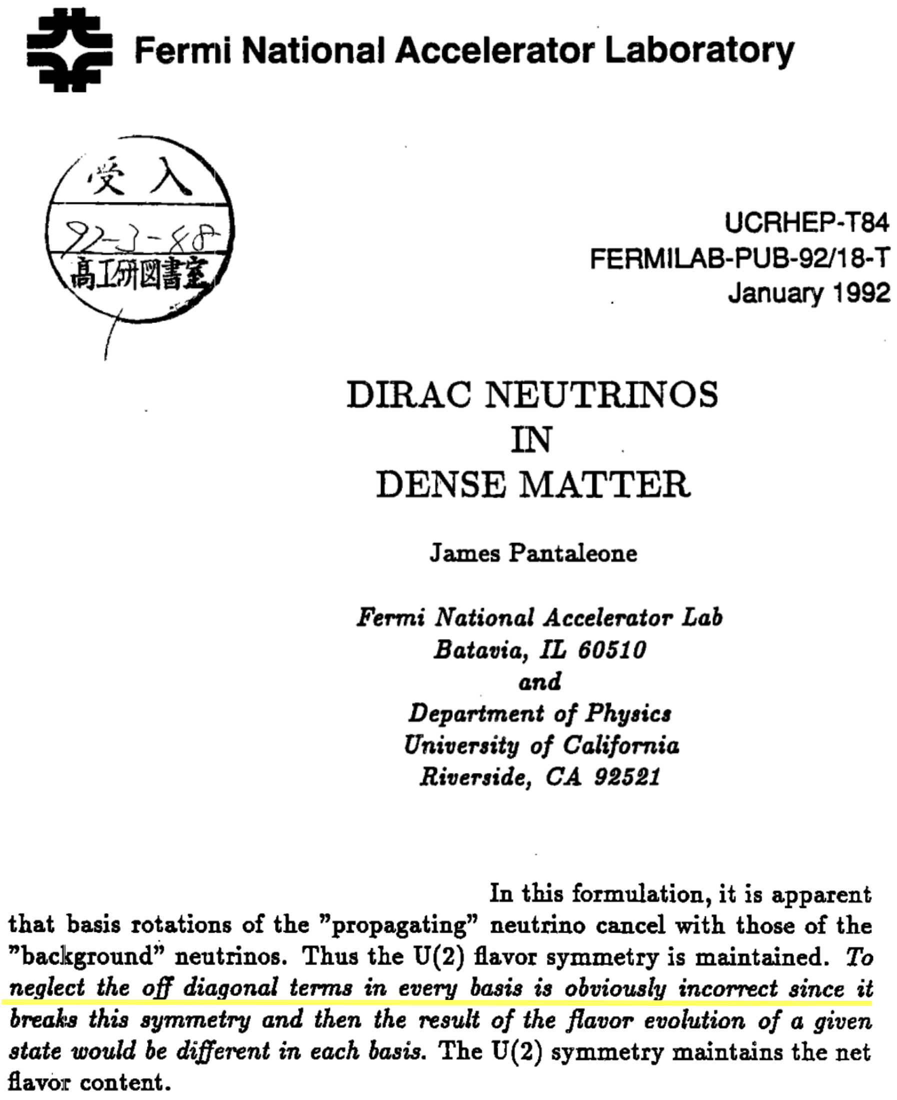
- DLS will detect electron neutrinos (+ all other flavors)
- Approximately 6% precision in survival probability with 1 kton-yr
- Can be a stringent test of the LMA-MSW solution
- May be a way to discover non-standard interactions and other exotic effects




ν_e Survival Probability

Bhavesh Chauhan + DLS Study Group

Collective Effects : Physical Origin



Fermi National Accelerator Laboratory

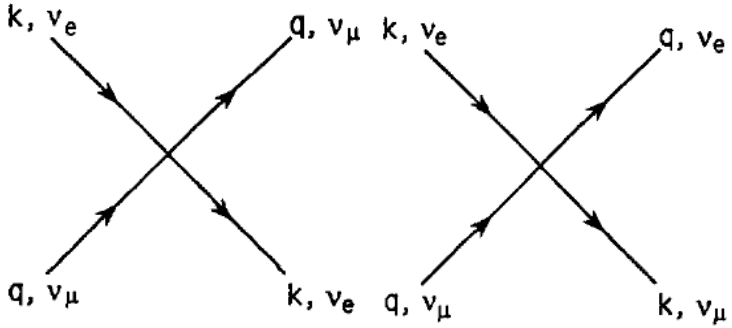
UCRHEP-T84
FERMILAB-PUB-92/18-T
January 1992

DIRAC NEUTRINOS
IN
DENSE MATTER

James Pantaleone

Fermi National Accelerator Lab
Batavia, IL 60510
and
Department of Physics
University of California
Riverside, CA 92521

In this formulation, it is apparent that basis rotations of the "propagating" neutrino cancel with those of the "background" neutrinos. Thus the U(2) flavor symmetry is maintained. *To neglect the off diagonal terms in every basis is obviously incorrect since it breaks this symmetry and then the result of the flavor evolution of a given state would be different in each basis.* The U(2) symmetry maintains the net flavor content.



Pantaleone (1992)

Forward scattering neutrinos can exchange flavor

$$\mathrm{i}\frac{\mathrm{d}}{\mathrm{d}t}\begin{vmatrix} |\nu_e(\boldsymbol{k})\nu_e(\boldsymbol{q})\rangle \\ |\nu_e(\boldsymbol{k})\nu_\mu(\boldsymbol{q})\rangle \\ |\nu_\mu(\boldsymbol{k})\nu_e(\boldsymbol{q})\rangle \\ |\nu_\mu(\boldsymbol{k})\nu_\mu(\boldsymbol{q})\rangle \end{vmatrix} = V_2 \begin{vmatrix} |\nu_e(\boldsymbol{k})\nu_e(\boldsymbol{q})\rangle \\ |\nu_e(\boldsymbol{k})\nu_\mu(\boldsymbol{q})\rangle \\ |\nu_\mu(\boldsymbol{k})\nu_e(\boldsymbol{q})\rangle \\ |\nu_\mu(\boldsymbol{k})\nu_\mu(\boldsymbol{q})\rangle \end{vmatrix}$$

where $V_{2\nu} = \sqrt{2}G_{\mathrm{F}}\xi\frac{1}{V}\begin{vmatrix} 2 & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 2 \end{vmatrix}$

The Basic Idea

$$i\partial_t|\nu_i\rangle = \left(\sum_{j=1}^{N}(1-\hat{p}_j\cdot\hat{p}_i)\underbrace{|\nu_j\rangle\langle\nu_j|}_{\rho_j} + \ldots\right)|\nu_i\rangle$$

Collective effect

Usual terms

Neutrinos give a phase-shift to other neutrinos.

This couples the linear equations and makes them “nonlinear”.

Can linearize and ask -- is the “evolution frequency” complex? If yes, Instability.

SN Neutrino Oscillations

$$i\left(\partial_t + \mathbf{v}_\mathbf{p}\cdot\partial_\mathbf{p}\right)\rho_\mathbf{p} = +\left[\frac{\mathsf{M}^2}{2E},\,\rho_\mathbf{p}\right] + \sqrt{2}G_\mathrm{F}\left[\mathsf{L},\rho_\mathbf{p}\right] + \sqrt{2}G_\mathrm{F}\int\frac{d^3\mathbf{q}}{(2\pi)^3}\left(1-\hat{\mathbf{p}}\cdot\hat{\mathbf{q}}\right)\left[\rho_\mathbf{q}-\bar{\rho}_\mathbf{q},\rho_\mathbf{p}\right]$$

Vacuum oscillations depend on neutrino mass matrix M
Overall minus sign for antineutrinos

MSW effect depends on ordinary matter density L, i.e. mainly electron density

Collective effects depends on the neutrino density

$$\omega = \frac{\Delta m^2}{2E}$$

$$\lambda = \sqrt{2}G_F\, n_e$$

$$\mu = \sqrt{2}G_F\, n_\nu$$

In general, a 7 dimensional problem
3 momentum (E, θ_p, ϕ_p) + 3 space (r, θ, ϕ) + 1 time (t)
Dimensionality of calculations denoted by $n_p + n_x + n_t$

Collective Oscillation: Instability

When density is high

$$\mu = G_F n_\nu \gg \frac{\Delta m^2}{2E} = \omega$$

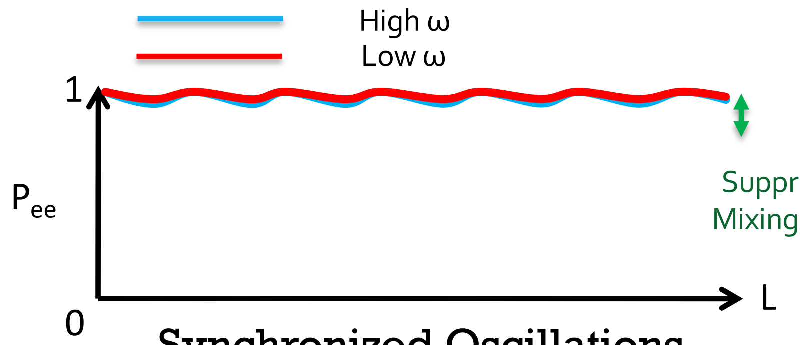
the collective oscillations have small amplitude

As neutrino density gets lower

$$G_F\, n_\nu \gtrsim \frac{\Delta m^2}{2E}$$

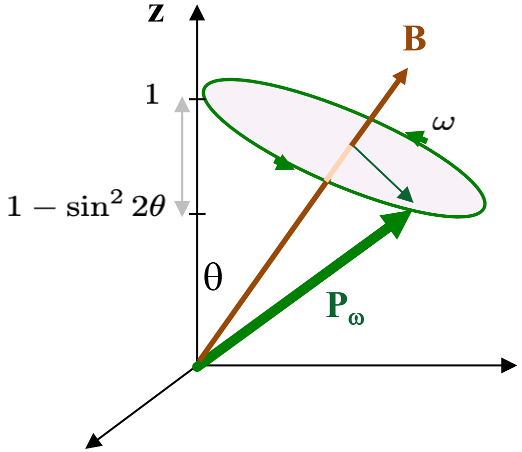
the system can be unstable

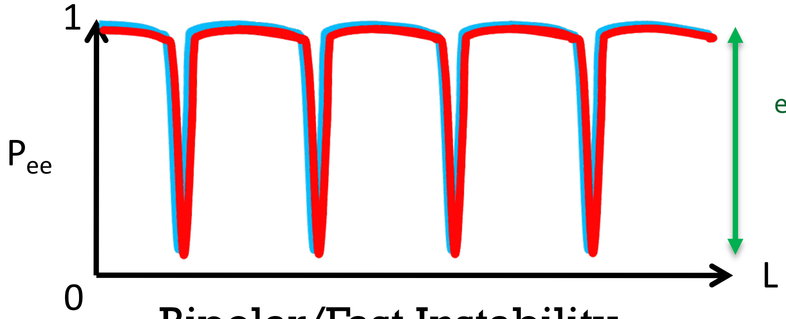
Instability grows at rate sqrt(ωμ) … (slow)
or proportional to μ … (fast)



Suppressed Mixing Angle

Synchronized Oscillations





Large Change even with Small Mixing Angle

Bipolar/Fast Instability

Pioneering work by Pantaleone, Kostelecky, Samuel in the 90s
Second-wave in 2005-06 by Duan, Fuller, Carlson, Qian + others
Fast Oscillations by Sawyer (2015)
Chakraborty, Hansen, Izaguirre, Raffelt (2016)
Dasgupta, Mirizzi, Sen (2016)

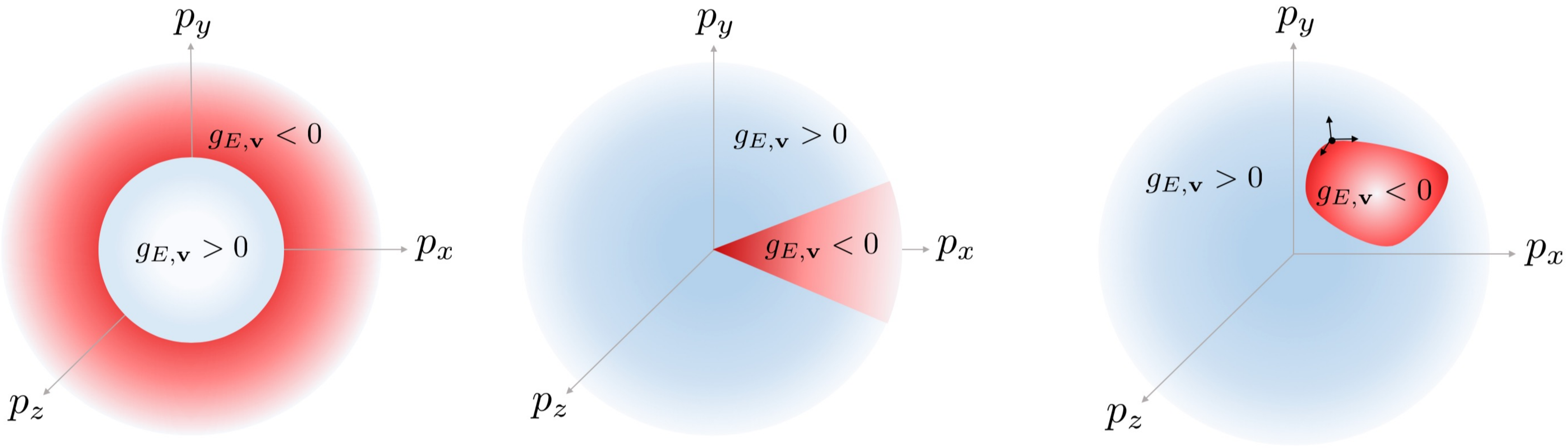
Crossing Theorem

Collective instability occurs *only if* momentum distributions of any two flavors *cross each other* around some momentum

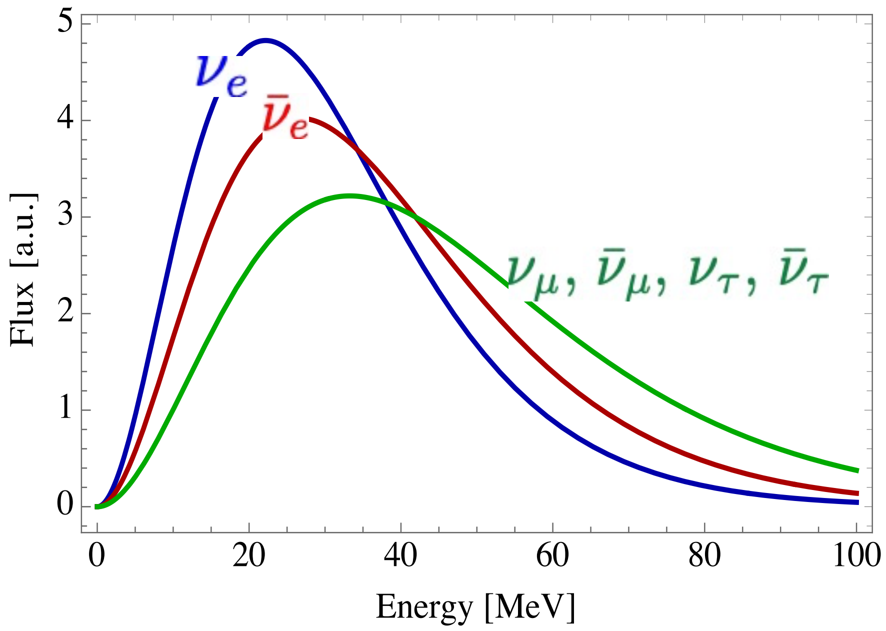
This is related to the positive-definiteness of a matrix

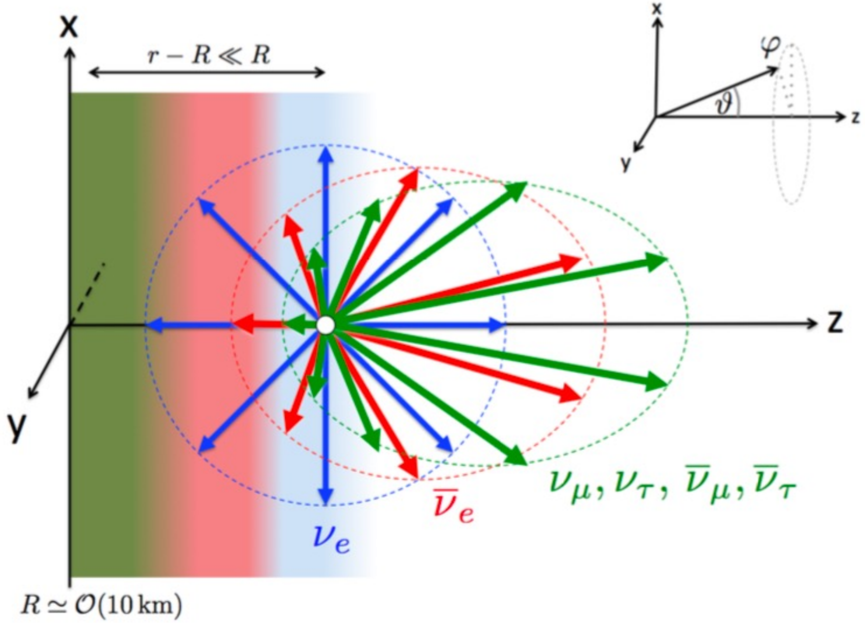
Dasgupta (2110.00192 ; PRL 2022)
Morinaga (PRD, 2022)

$$g_\Gamma = \sqrt{2} G_\mathrm{F} \begin{cases} f_{\nu_e,\mathbf{p}} - f_{\nu_\mu,\mathbf{p}} & \text{for } E > 0 \\ f_{\bar{\nu}_\mu,\mathbf{p}} - f_{\bar{\nu}_e,\mathbf{p}} & \text{for } E < 0 \end{cases}$$



Why crossings exist





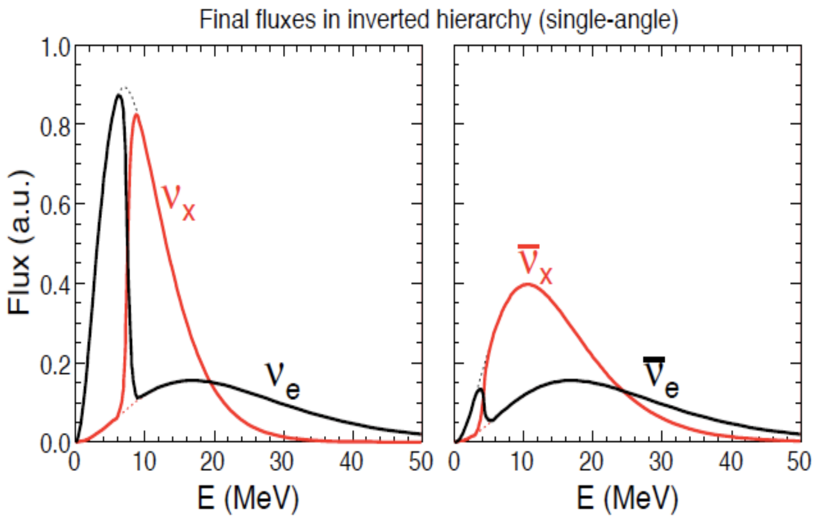
Different flavors have different energy spectrum

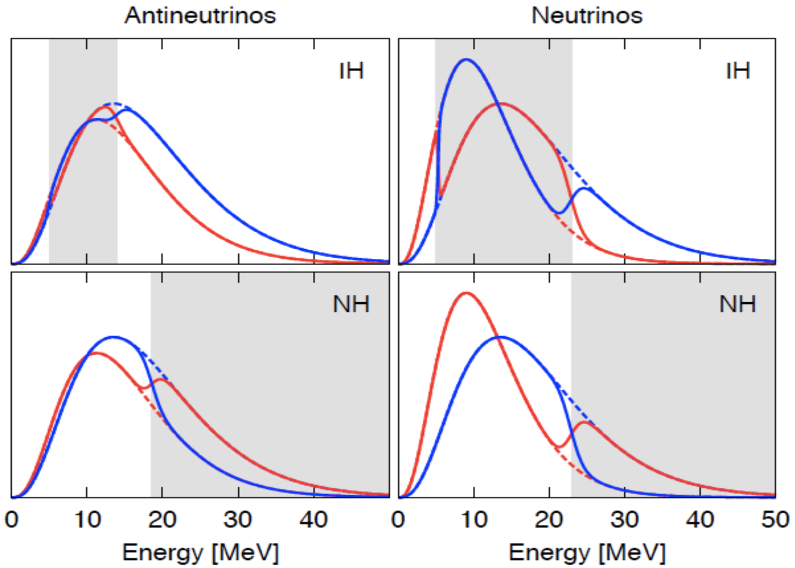
Crossing leads to (slow) instability

Different flavors have different angular distribution

Crossing leads to (fast) instability

Spectral Swaps due to Slow Effects





Portions of the energy spectra get exchanged

Initially thought to occur for Inverted ordering

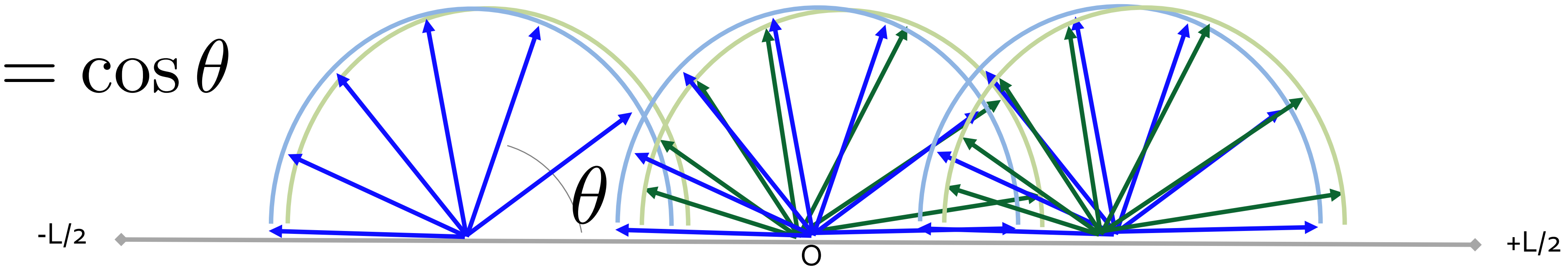
Later realized that this occurs for both orderings and there can be multiple spectral splits

Seminal papers by Duan, Fuller, Carlson, Qian (2005, 2006, 2007)
Raffelt and Smirnov (2007, 2007)
Fogli, Lisi, Marrone, Mirizzi (2008)

Dasgupta, Dighe, Raffelt, Smirnov (2009)
Friedland (2010)

Fast Oscillation : Numerics

$$v = \cos\theta$$



1 dimensional box of length L

$$(\partial_t + v\partial_z)\mathsf{S}_v = \mu_0 \int_{-1}^{+1} dv' G_{v'}\,(1 - vv')\,\mathsf{S}_{v'} \times \mathsf{S}_v$$

Neutrinos labeled by velocity emitted in electron or muon state at a small region of size L

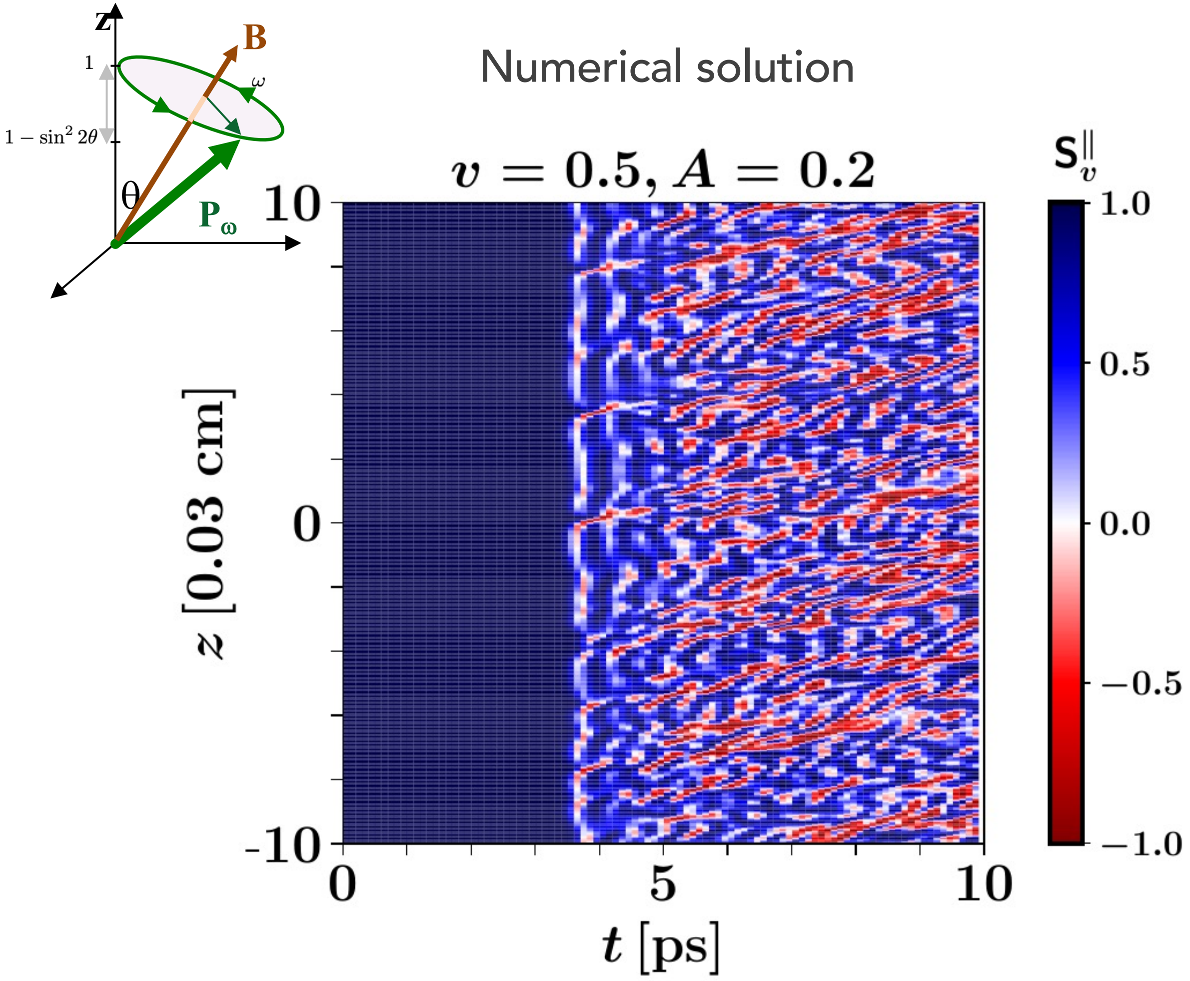
Net-emission at velocity v is $G_v = f_e(v) - f_{muon}(v)$

which is the difference of phase space distributions of the two flavors at each v.

1+1+1d calculation Bhattacharyya and Dasgupta (PRD 2020 ; PRL 2021, 2205.05129)

Fast Mixing due to Fast Instability

Numerical solution



Bhattacharyya and Dasgupta (PRD, 2020 ; PRL 2021)

Analytical Understanding via Coarse-graining

$$\partial_t \langle M_n \rangle = \frac{\langle M_1 \rangle}{2} \left(\partial_n^2 \langle M_n \rangle + \frac{1}{n} \partial_n \langle M_n \rangle \right)$$

A diffusion of the "difference of flavors" to higher multipoles of emission angle (i.e., momentum)

Several other groups have since obtained similar results

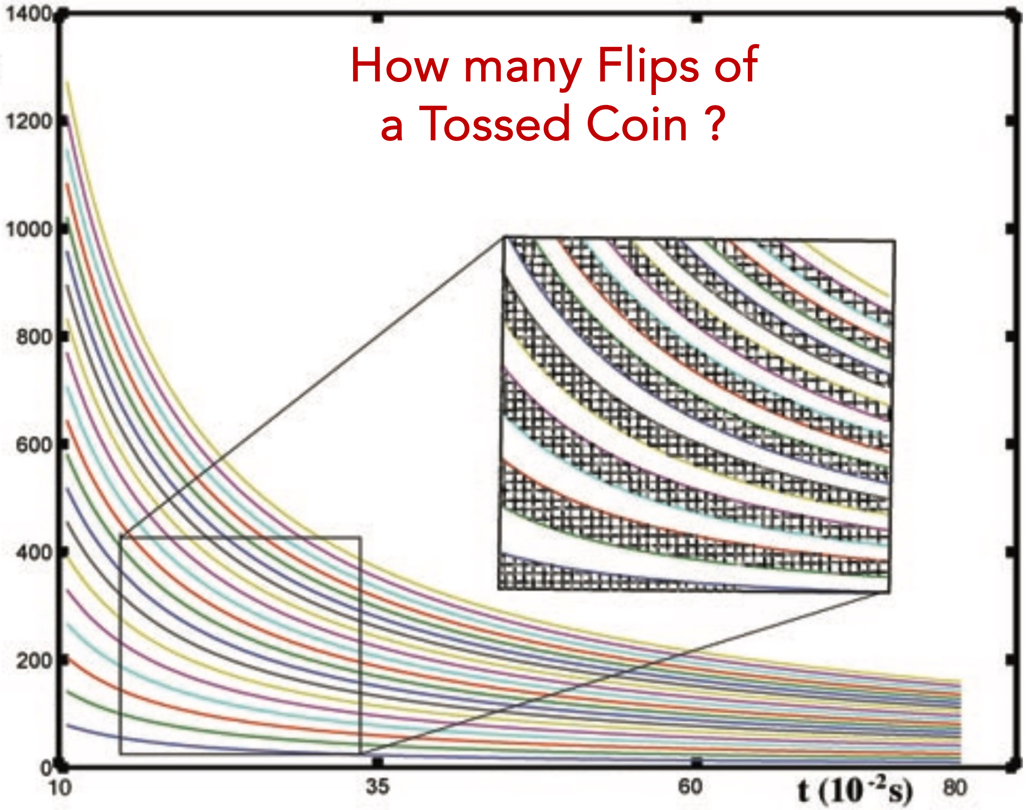
- Richers et al. @ Berkeley
- Wu et al. @ Taiwan
- Sigl @ Hamburg

Survival Probability starts at 1

Oscillates coherently a few times

And then decoheres

Probability from Certainty



How many Flips of a Tossed Coin ?

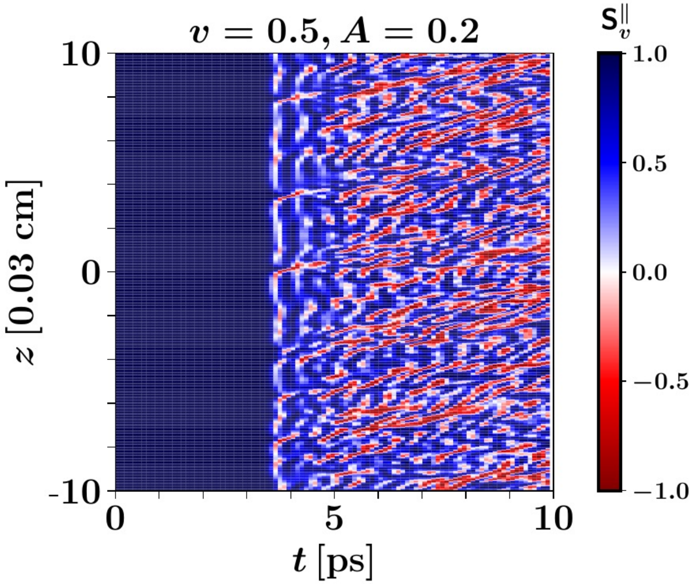
Initial Angular Velocity

Time in air ~ Initial Velocity

Fig. from Diaconis, Holmes, Mongomery (2007)

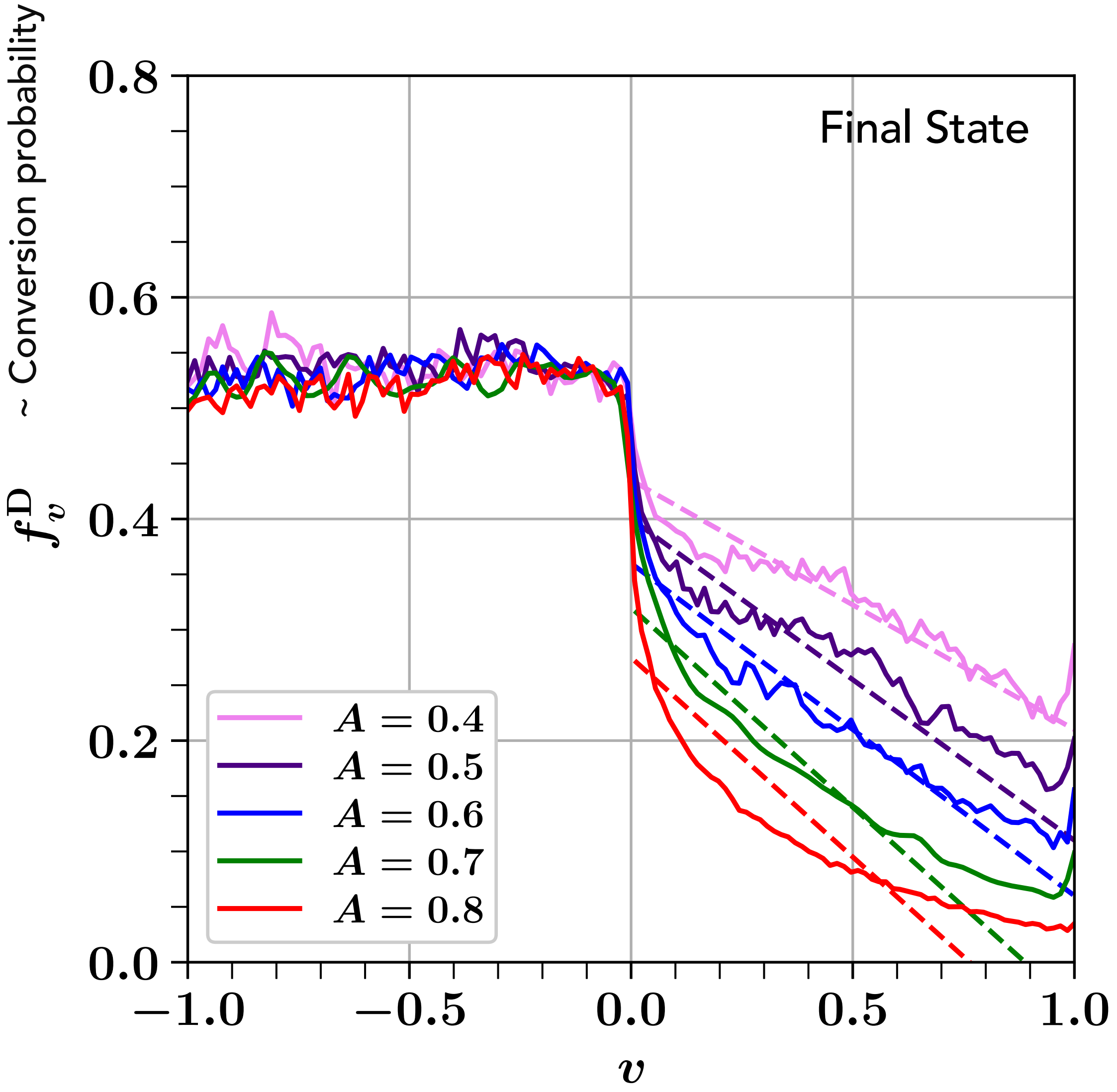
Coarse-grained measurements lead to probabilities via strong averaging over deterministic possibilities

Coarse-graining can lead to apparent irreversibility even if the evolution is reversible



Coarse-grained probability

Fast Depolarization



$$f_v^{\mathrm{D}} \approx \begin{cases} \frac{1}{2} - \frac{A}{4} - \frac{3A}{8}\,v, & \text{if } v > 0 \\ \frac{1}{2}, & \text{if } v < 0 \end{cases}$$

Bhattacharyya and Dasgupta (PRL, 2021)

Can approximately predict the spatially coarse-grained degree of flavor mixing *without* full numerical solution.

If you do not know how to include collective oscillations in your SN neutrino analysis what should you do?

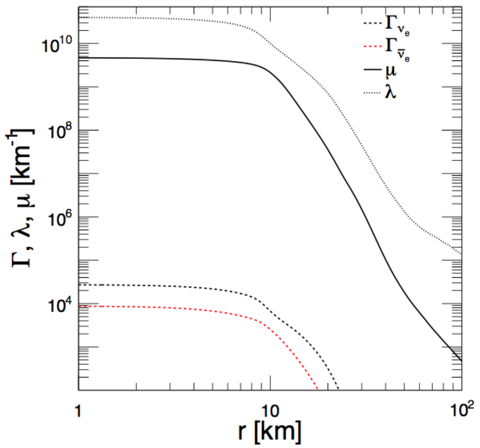
Assume "Flavor Equilibrium"

if Crossing Exists

Bhattacharyya and Dasgupta (PRL, 2021)

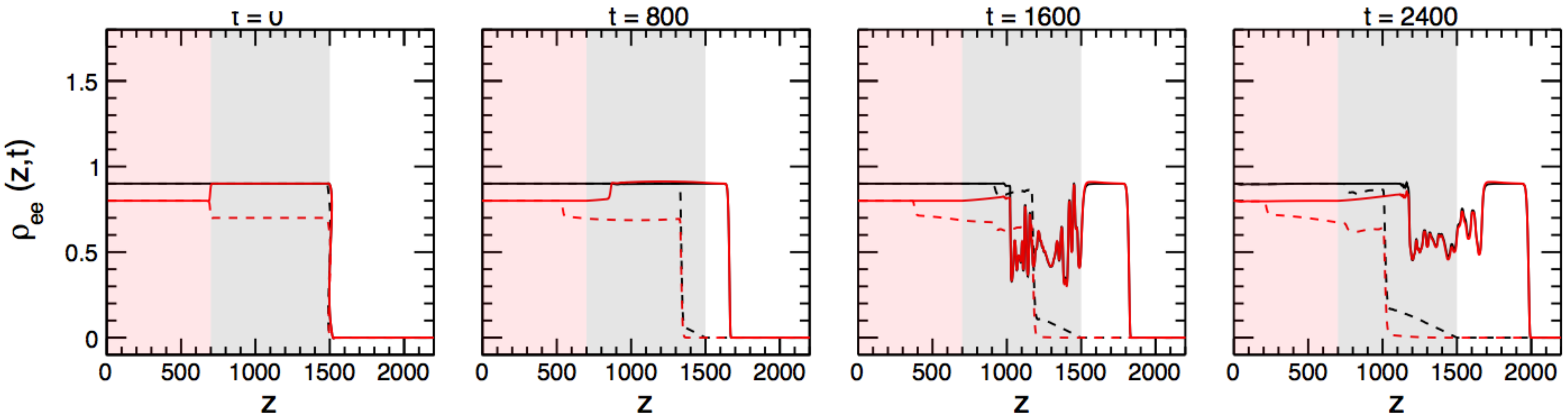
Caution: This is very simplified advice but a good starting point
One should do better : Partial Flavor Equilibrium
Need more realism: Contact your favorite SN Neutrino Physicist

Collisions vs. Oscillations



Collisions can be strong enough to create a difference between electron neutrinos and antineutrinos.
Yet, may not damp oscillations.

Capozzi, Dasgupta, Mirizzi, Sen (2018)
See also Johns (2021)
Martin, Carlson, Cirigliano, Duan (2021)
Sasaki, Takiwaki (2021)
Shalgar and Tamborra (2021)



Do fast conversions, once generated, penetrate the SN core?

Related work

- Mean-field treatment via Wigner fns.

Birol, Pehlivan, Balantekin, Kajino (2018)
Stirner, Sigl, Raffelt (2018)
Vlasenko, Fuller, Cirigliano (2014)
Volpe, Vaananen, Espinoza (2013)
Cardall (2008)

- Wavepackets/Kinematic decoherence

Akhmedov, Kopp, Lindner (2017)
Hansen, Smirnov (2016)

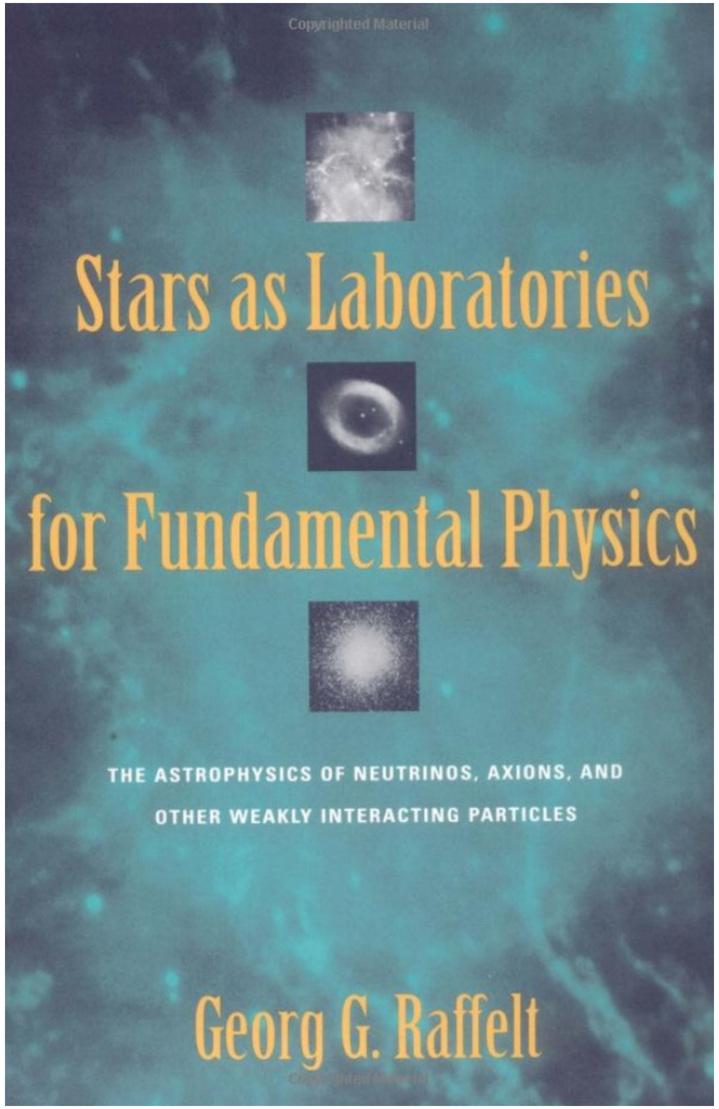
- Sterile nus, NSI, BSM, …

Skipping non-standard physics for lack of space and time

- Interpretation of collective effects

Hansen and Smirnov (2018)
Morinaga and Yamada (2017)

Fundamental Physics with Stars

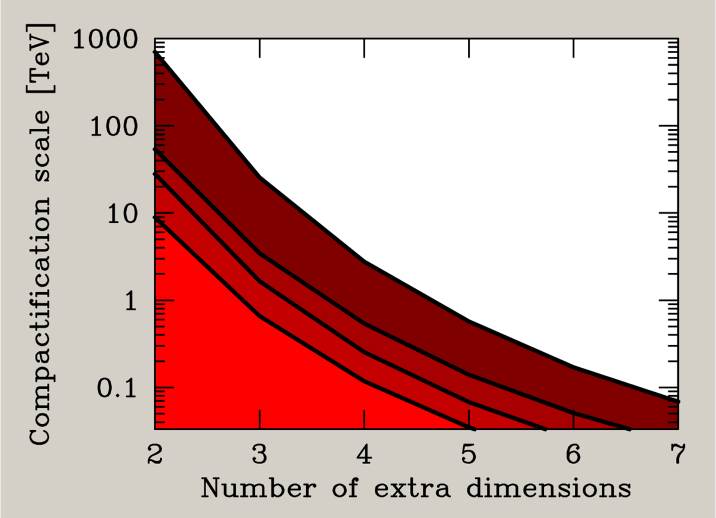


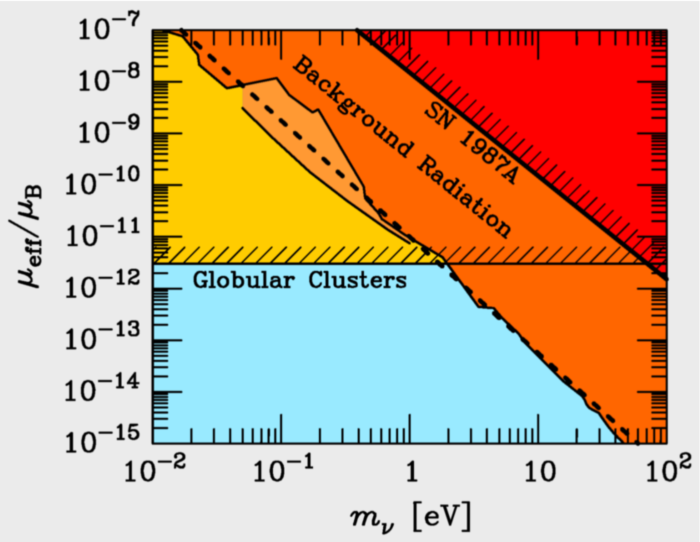
Stars as Laboratories for Fundamental Physics

THE ASTROPHYSICS OF NEUTRINOS, AXIONS, AND OTHER WEAKLY INTERACTING PARTICLES

Georg G. Raffelt

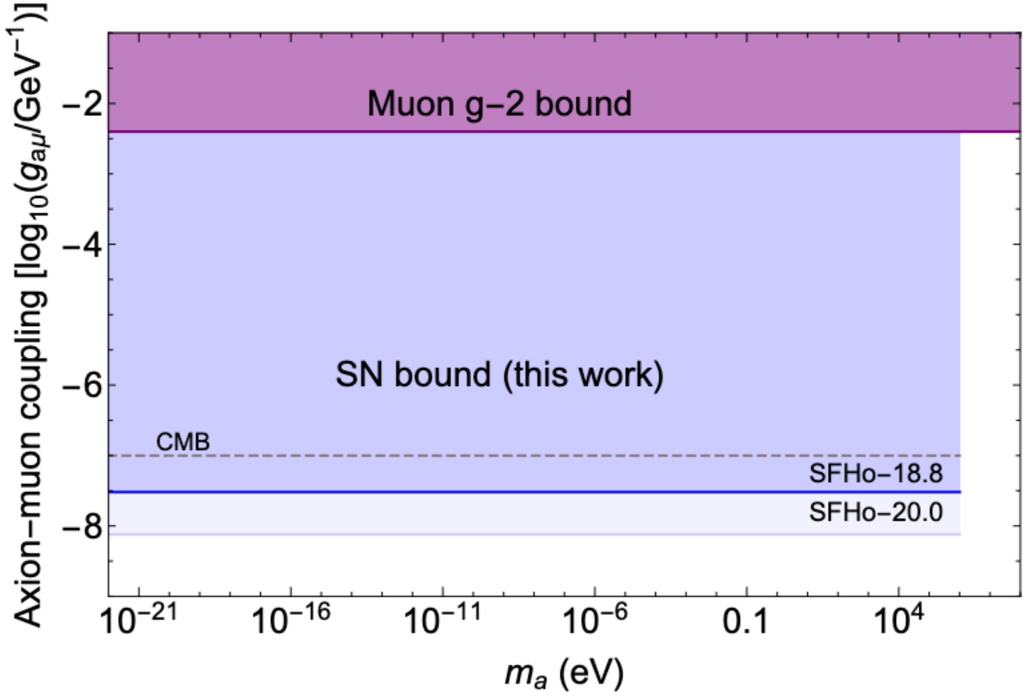
Copyright: The University of Chicago





New Ingredients: e.g. Muons

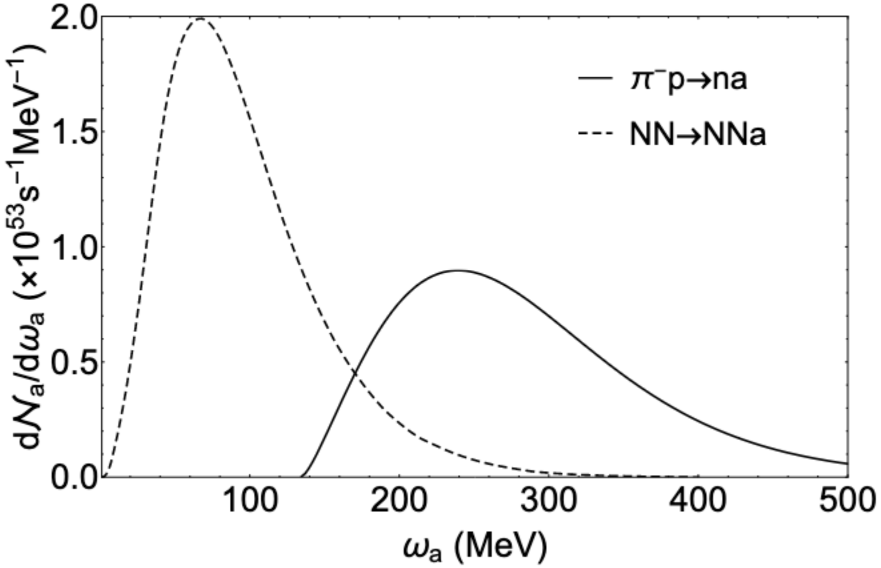
Temperature in SN ~ 30-60 MeV.
Muons are not heavily suppressed.
Can use this to put bounds on BSM physics coupling to muons, e.g., an axion coupling to muons, or a light feebly interacting $L_e - L_\mu$ boson.



Bollig et al. (PRL, 2017)
Bolling, de Rocco, Graham, Janka (PRL, 2017)
Croon, Elor, Leane, McDermott (2020)
Caputo, Raffelt, Vitagliano (PRD, 2022)

Better limits from better models of SN, and inclusion of previously ignored physical effects

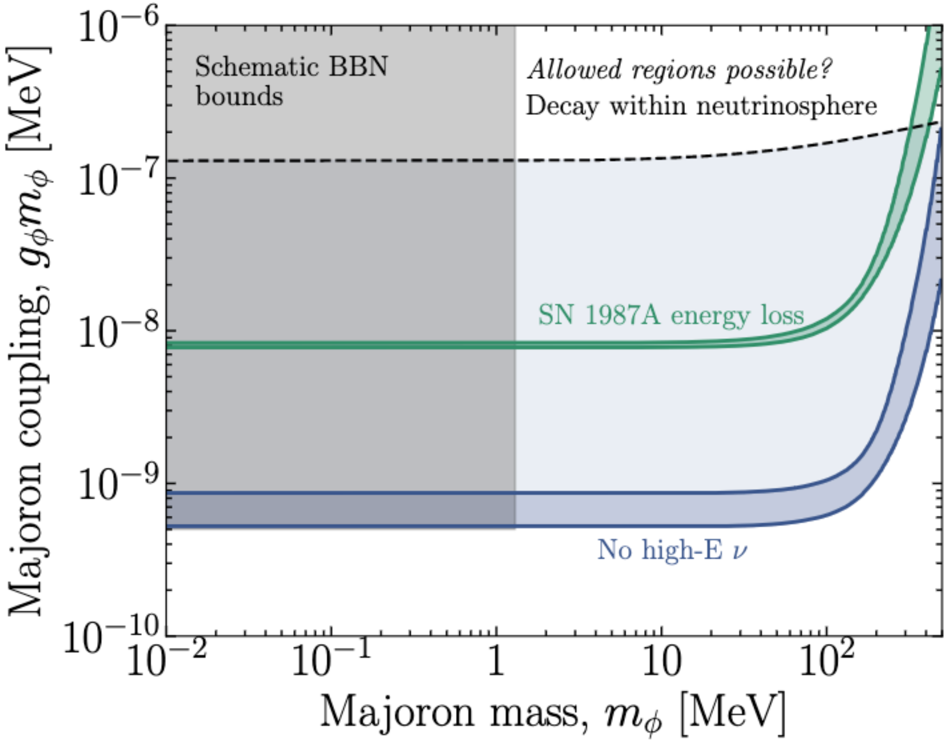
New Processes : e.g. NN -> NNa



Newly considered processes that produce axions more prominently at higher energies.

Fore, Reddy (PRC, 2020)
Carenza, Fore, Gianotti, Mirizzi, Reddy (PRL, 2020)

New Arguments: e.g. HENs

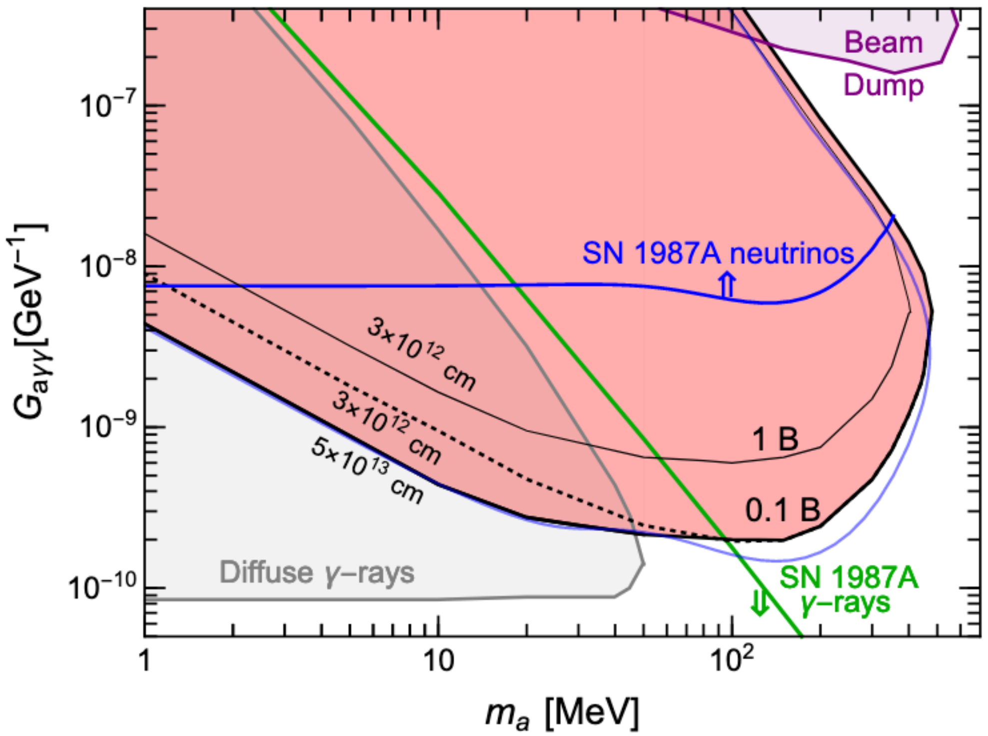


Fiorillo, Raffelt, Vitagliano (2209.11773)

A Majoron not only steals energy from core of SN, it can also decay back into high-E neutrinos.

Non-observation of high-E neutrinos from SN gives stronger bounds than cooling.

New Arguments: e.g. Ejecta Energy

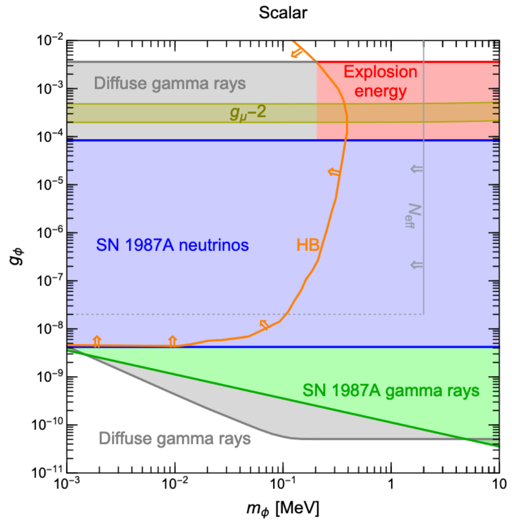


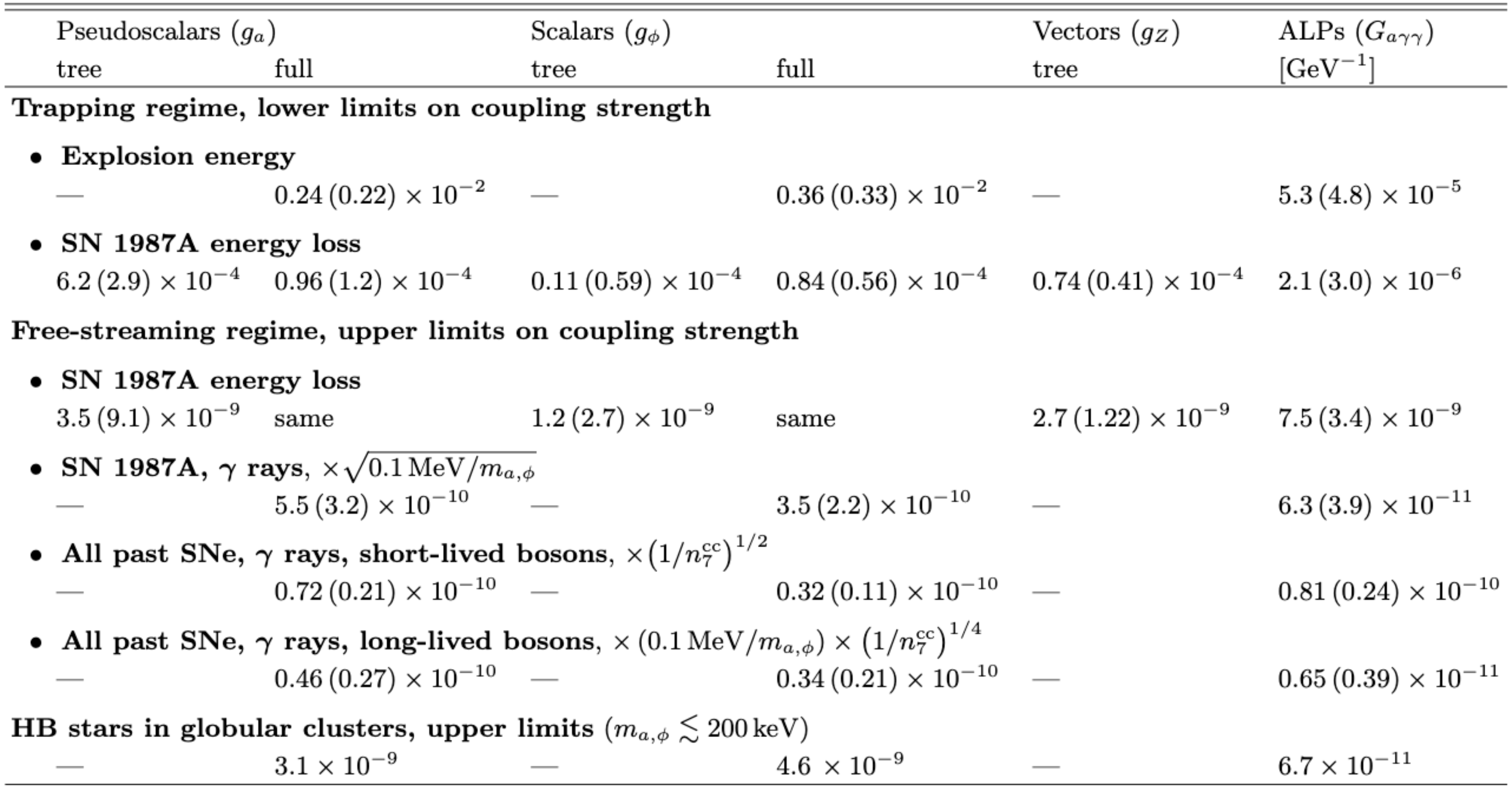
Caputo, Janka, Raffelt, Vitagliano (PRL, 2022)

A BSM particle produced in SN not only cools but also interacts and redeposits energy in ejecta

Observed ejecta energy can give stronger or complementary bounds.

SN Bounds Galore

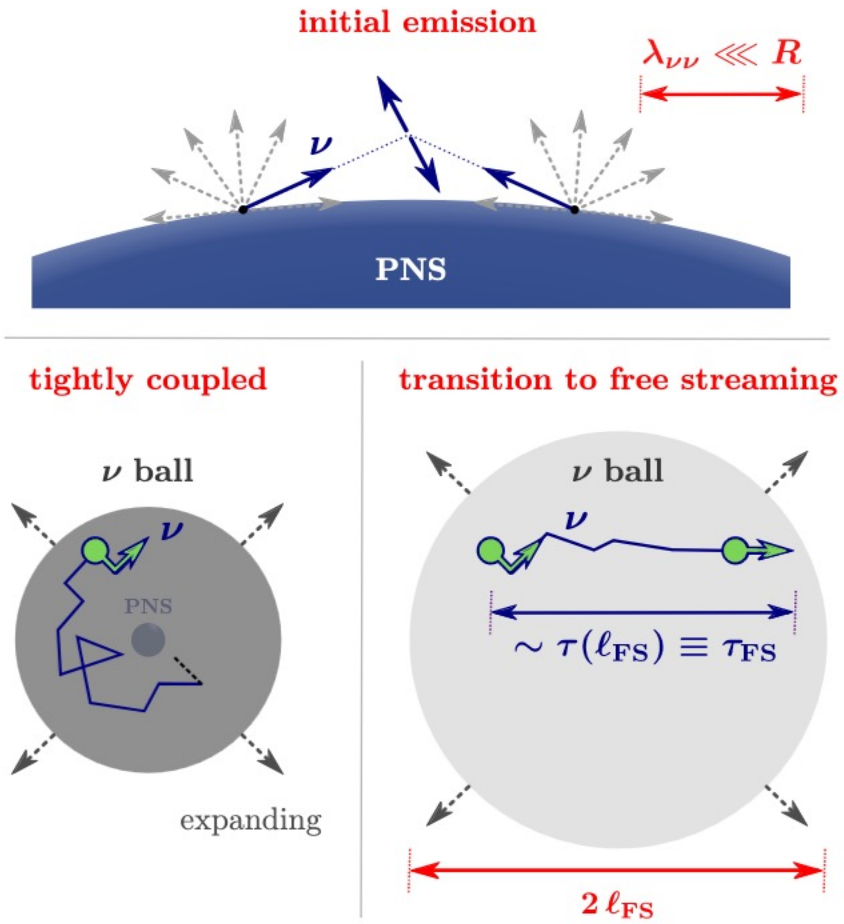




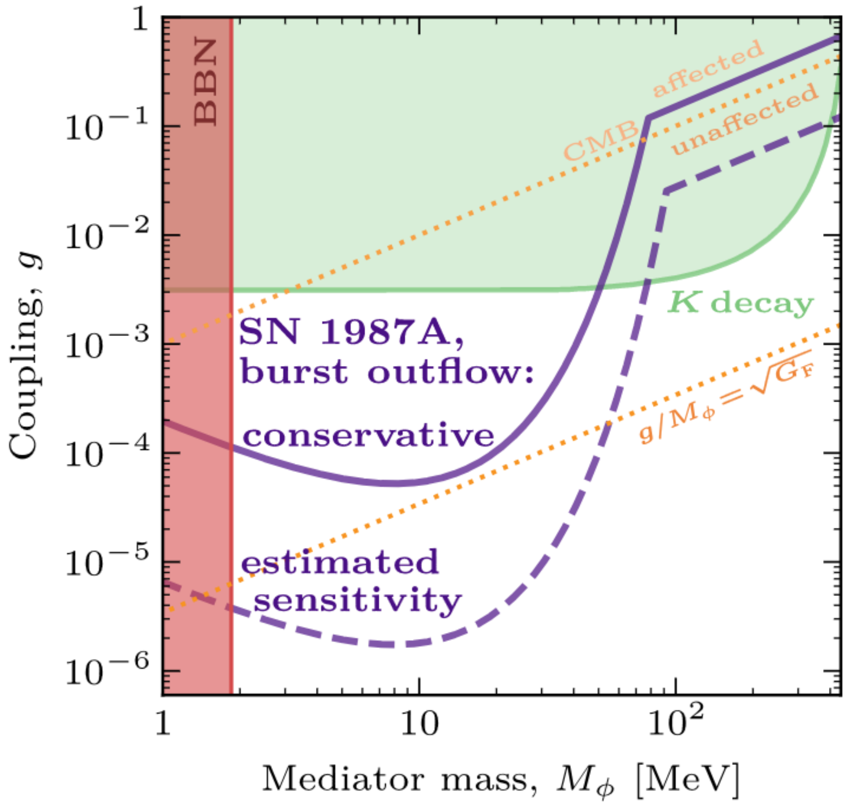
Pseudoscalars (g_a)		Scalars (g_ϕ)		Vectors (g_Z)	ALPs ($G_{a\gamma\gamma}$)
tree	full	tree	full	tree	[GeV^{-1}]
Trapping regime, lower limits on coupling strength					
• **Explosion energy**					
—	$0.24\,(0.22) \times 10^{-2}$	—	$0.36\,(0.33) \times 10^{-2}$	—	$5.3\,(4.8) \times 10^{-5}$
• **SN 1987A energy loss**					
$6.2\,(2.9) \times 10^{-4}$	$0.96\,(1.2) \times 10^{-4}$	$0.11\,(0.59) \times 10^{-4}$	$0.84\,(0.56) \times 10^{-4}$	$0.74\,(0.41) \times 10^{-4}$	$2.1\,(3.0) \times 10^{-6}$
Free-streaming regime, upper limits on coupling strength					
• **SN 1987A energy loss**					
$3.5\,(9.1) \times 10^{-9}$	same	$1.2\,(2.7) \times 10^{-9}$	same	$2.7\,(1.22) \times 10^{-9}$	$7.5\,(3.4) \times 10^{-9}$
• **SN 1987A, γ rays**, $\times\sqrt{0.1\,\mathrm{MeV}/m_{a,\phi}}$					
—	$5.5\,(3.2) \times 10^{-10}$	—	$3.5\,(2.2) \times 10^{-10}$	—	$6.3\,(3.9) \times 10^{-11}$
• **All past SNe, γ rays, short-lived bosons**, $\times\,(1/n_7^{\rm cc})^{1/2}$					
—	$0.72\,(0.21) \times 10^{-10}$	—	$0.32\,(0.11) \times 10^{-10}$	—	$0.81\,(0.24) \times 10^{-10}$
• **All past SNe, γ rays, long-lived bosons**, $\times\,(0.1\,\mathrm{MeV}/m_{a,\phi}) \times (1/n_7^{\rm cc})^{1/4}$					
—	$0.46\,(0.27) \times 10^{-10}$	—	$0.34\,(0.21) \times 10^{-10}$	—	$0.65\,(0.39) \times 10^{-11}$
HB stars in globular clusters, upper limits ($m_{a,\phi} \lesssim 200\,\mathrm{keV}$)					
—	3.1×10^{-9}	—	4.6×10^{-9}	—	6.7×10^{-11}

Caputo, Raffelt, Vitagliano (PRD, 2022)
Cerdeno, Cermeno, Farzan (2023)
+ many more papers … apologies for incomplete list

Puzzle: Neutrino Self-Interactions



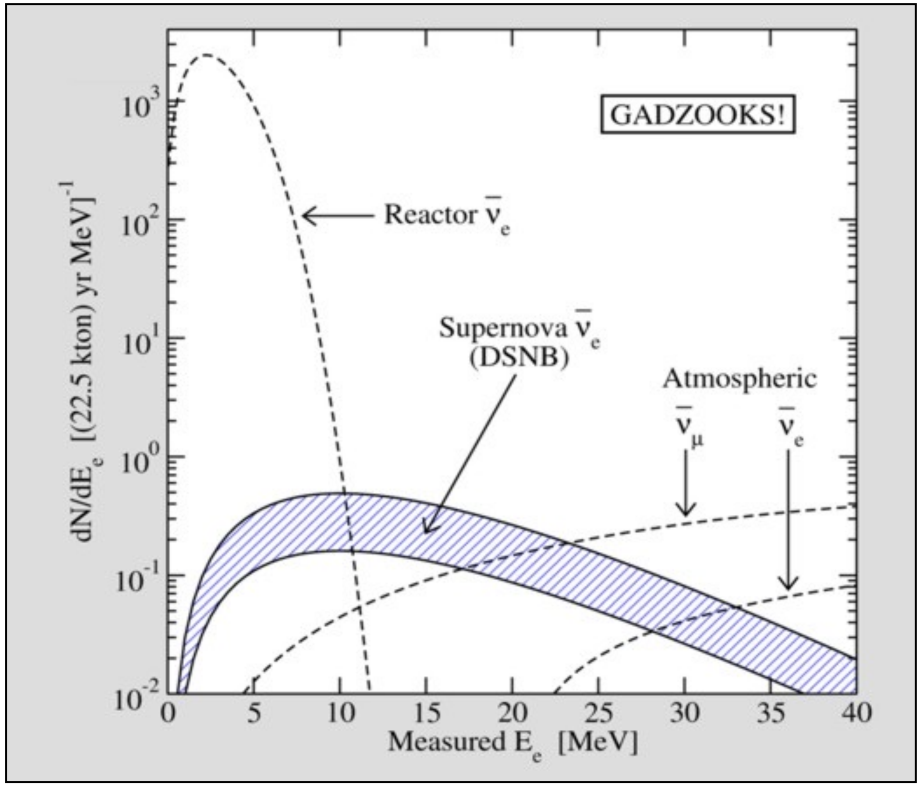
Chang, Isteban, Beacom, Thompson, Hirata (2206.12426)
+ previous papers …



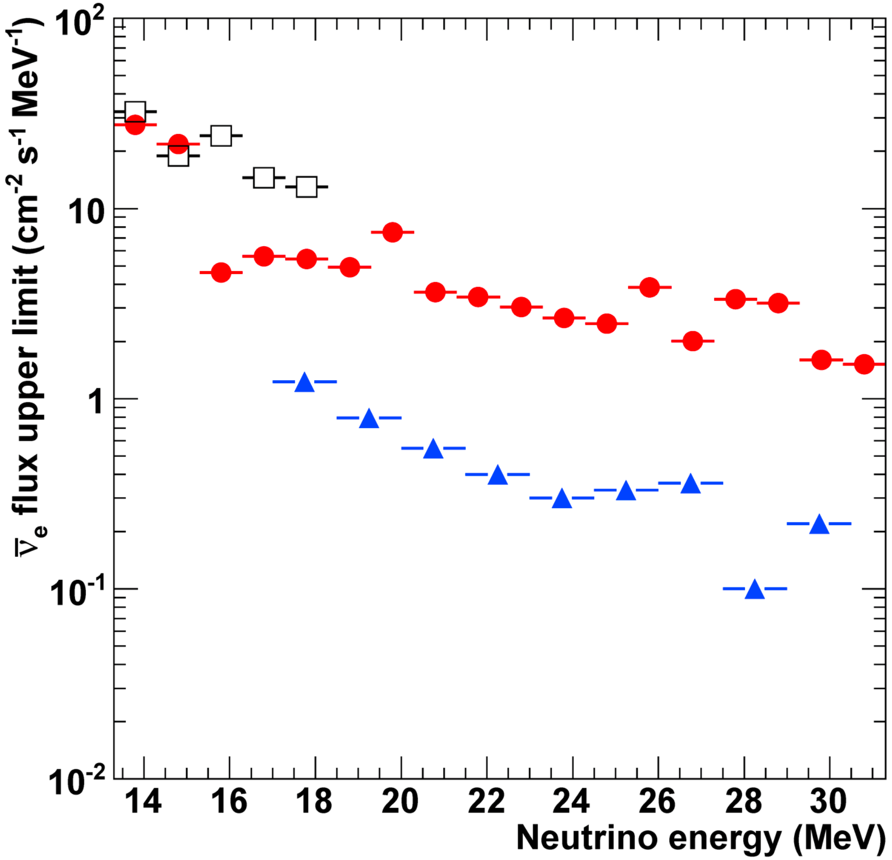
If neutrinos self-decouple too late
they behave like a fluid whose behavior
is still not fully understood

Many more …

DSNB



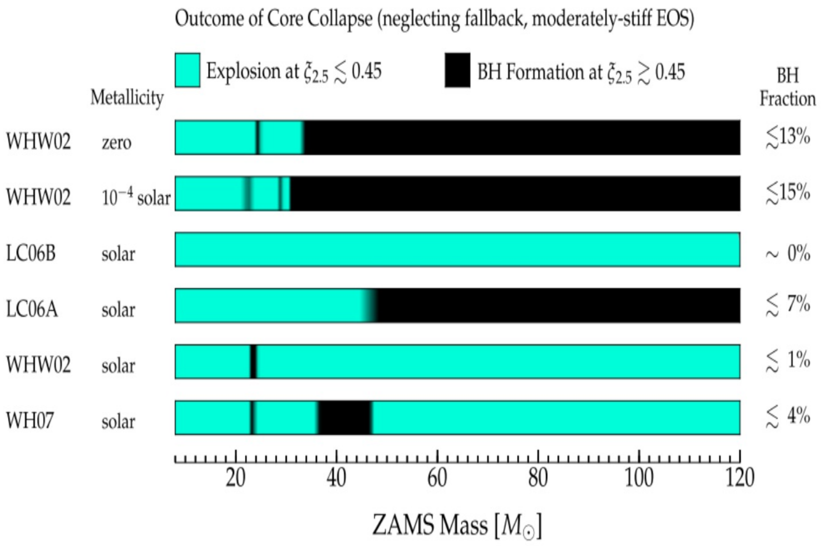
Zeldovic (1964), Bisnovatyi-Kogan and Seidov (1982)
Plot from Beacom and Vagins (2003)
See reviews by Lunardini (2010) and Beacom (2010)

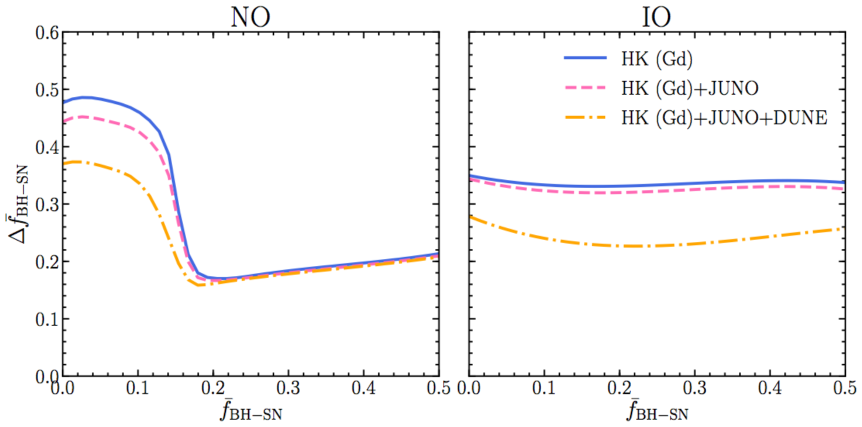


Super-Kamiokande Collaboration (2013)

We may be close to detecting the DSNB.
With Gd upgrade of Super-K this can be very promising.

Are there failed SN?





Some supernovae are not expected to explode

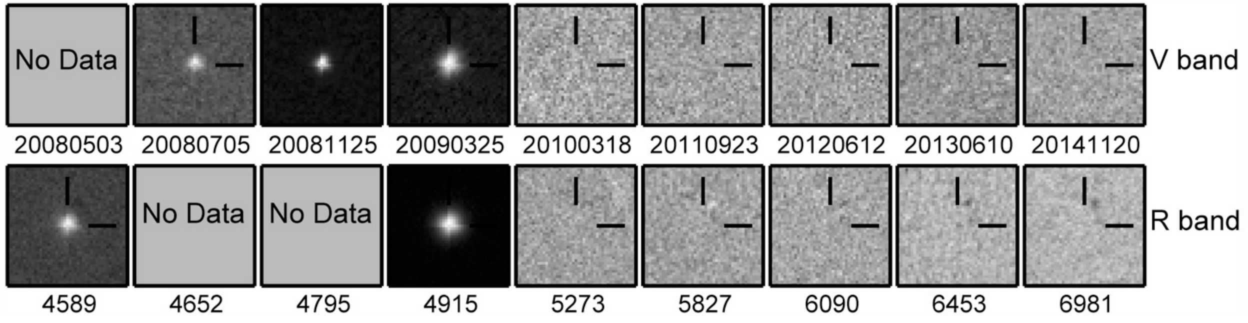
O'Connor and Ott (2013)

Ertl, Janka, Woosley, Sukhbold, Ugliano (2016)

DSNB is sensitive to the failed SN fraction

Lunardini (2009)

Moller, Suliga, Tamborra, Denton (2018)

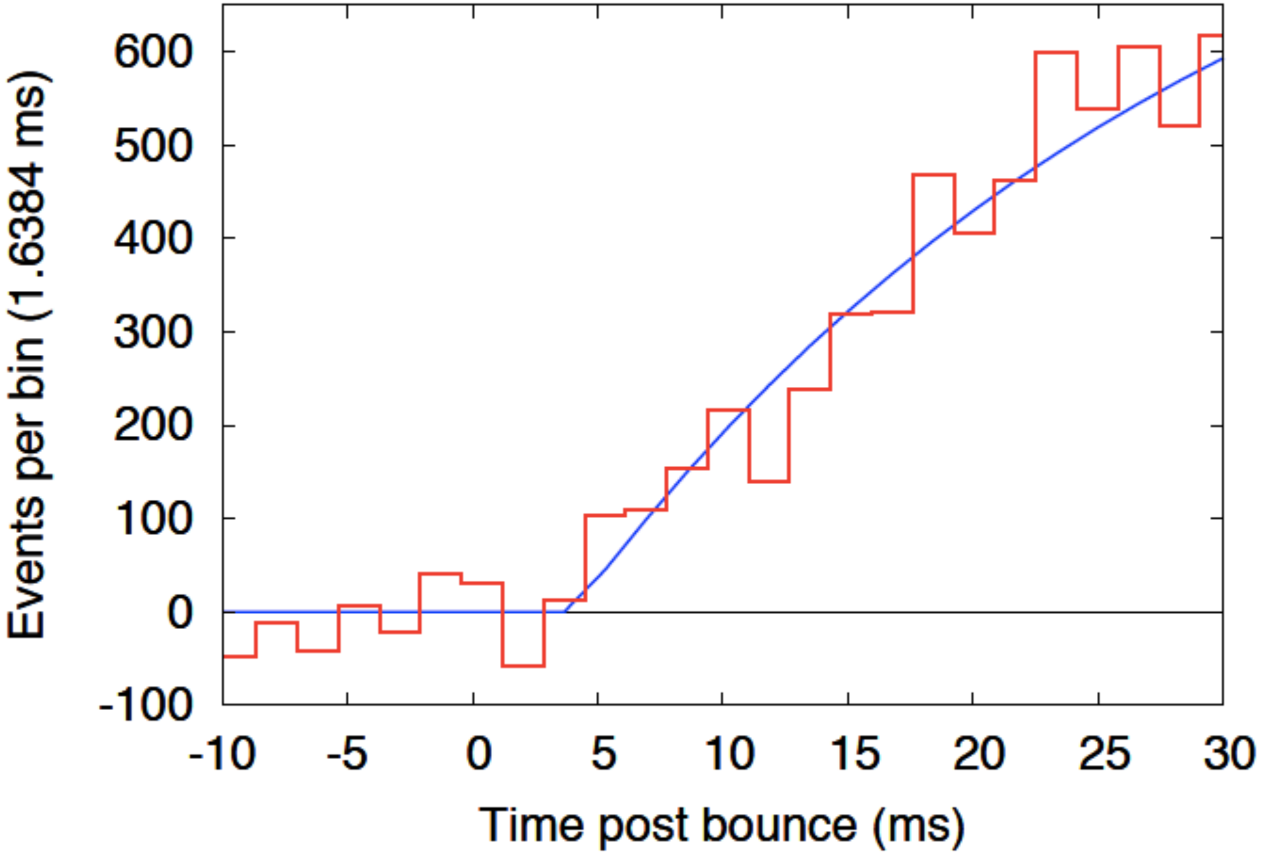


SN without a bang

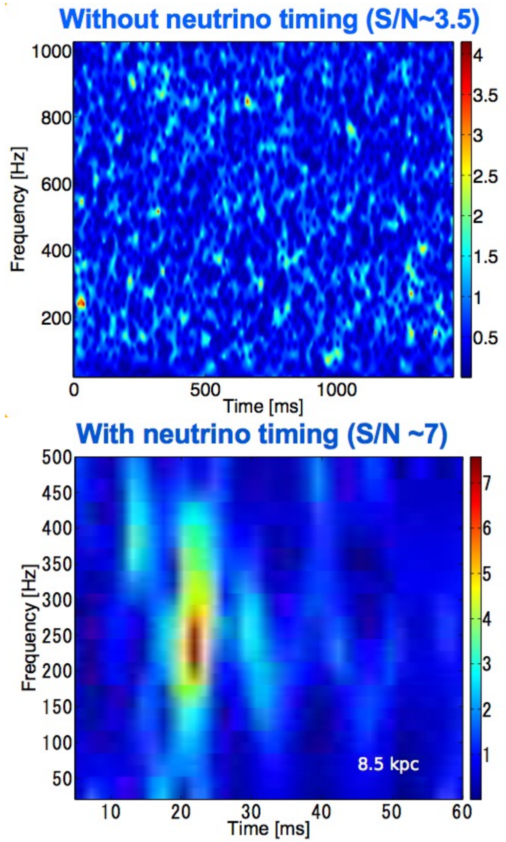
Gerke, Kochanek, Stanek (2014)

Reynolds, Fraser, Gilmore (2015)

Timing



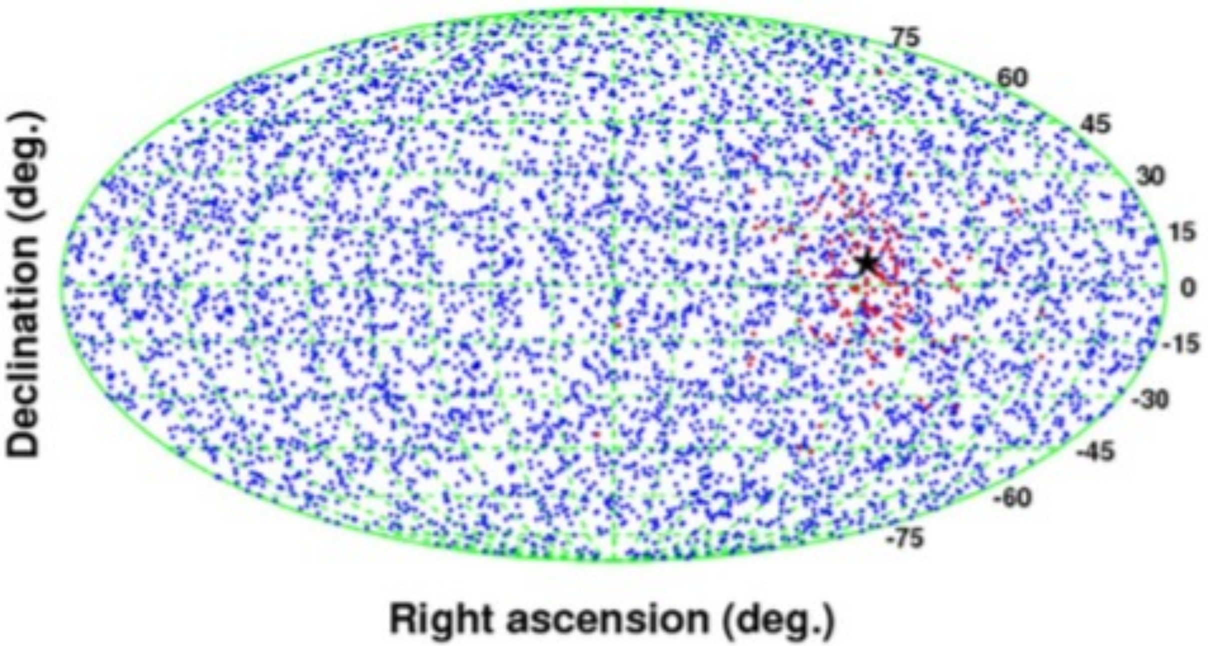
Pagliaroni, Vissani, Coccia, Fulgione (2009)
Plot from Halzen and Raffelt (2009)



Nakamura, Horiuchi, Tanaka, Hayama, Takiwaki, Kotake (2016)

Improved ability to spot the signal with different messengers

Pointing



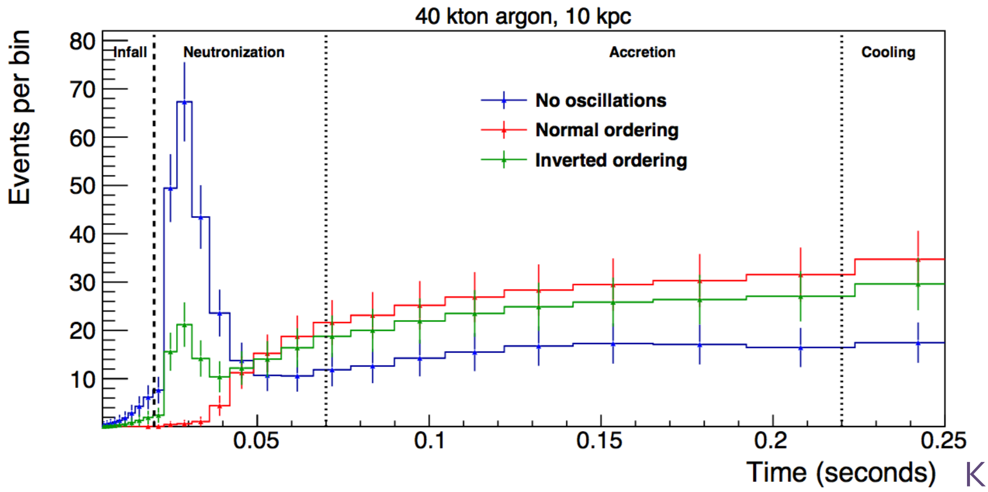
Using directionality of elastic scattering events and subtraction of tagged inverse beta "background"

Beacom and Vogel (1998)
Beacom and Vagins (2000)
Tomas, Semikox, Raffelt, Kachelriess, Dighe (2003)
Plot from Abe at al. for Super-K (2016)

For triangulation: see Beacom and Vogel (1998)
Muhlbeier, Nunokawa, and Zukanovich Funchal (2013)
Brdar, Lindner, Xu (2018)

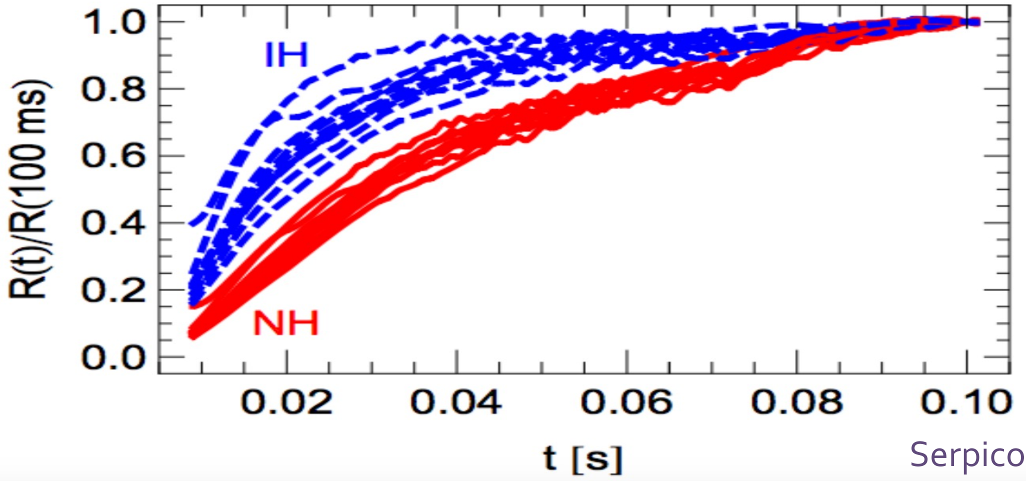
Pointing accuracy of a few degrees for SN at 10 kpc

Mass Ordering



No neutronization peak seen in electron neutrinos for inverted mass ordering

Plot from E. Worcester's talk at Neutrino 2018
Wallace, Burrows, Dolence (2015)
Kachelriess, Tomas, Buras, Janka, Marek, Rampp (2004)



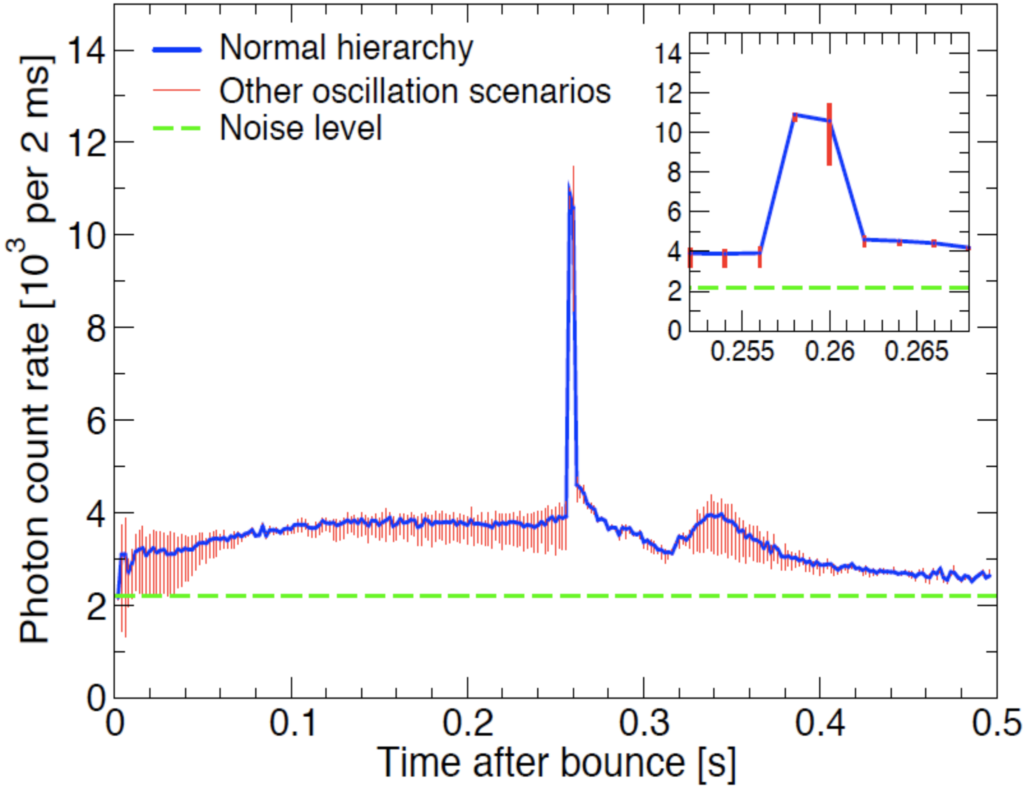
Electron antineutrino signal rises faster for inverted ordering

Note: This can change if fast conversions occurs in the accretion phase

Serpico, Chakraborty, Fischer, Hudepohl, Janka, Mirizzi (2011)

Neutronization burst can reveal the neutrino mass ordering

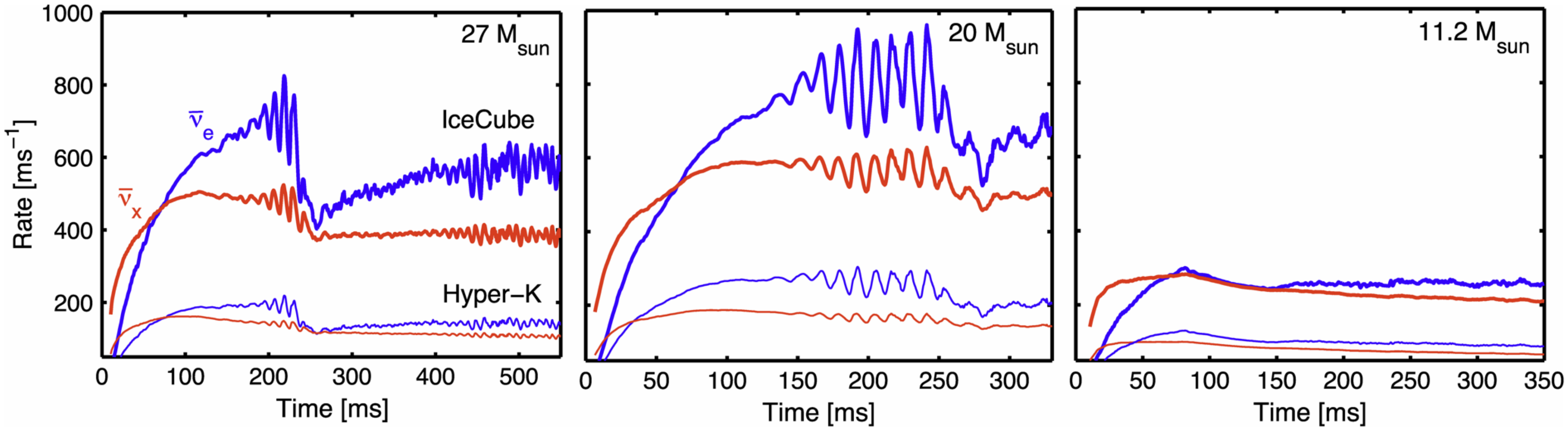
QCD transition in SN



Dasgupta, Fischer, Horiuchi, Liebendoerfer, Mirizzi, Sagert, Schaffner-Bielich (2009)
Simulation by Sagert, Fischer, Hempel, Pagliara, Schaffner-Bielich, Mezzacappa, Thielemann, Liebendoerfer (2008)

Neutrino telescopes are exquisitely sensitive to anything
that affects the electron antineutrino lightcurve

SASI Signatures



27 M_{sun} star
Multiple SASI episodes and convection

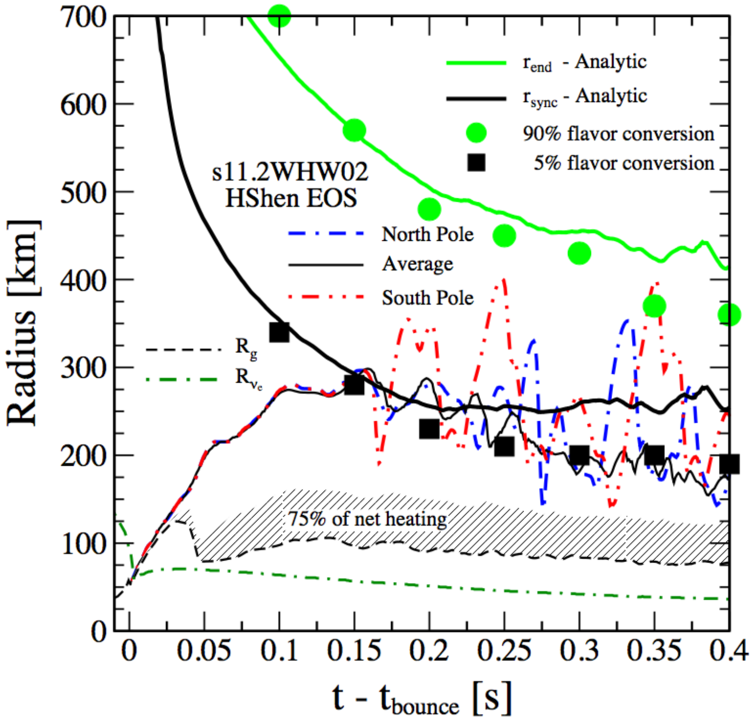
20 M_{sun} star
Single SASI episode and convection

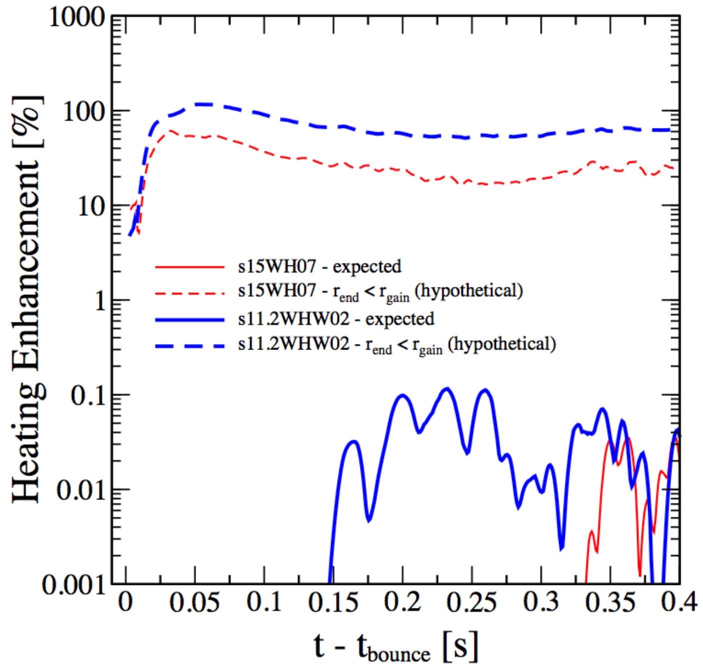
11.2 M_{sun} star
No SASI episode; only convection

Tamborra, Raffelt,Hanke, Janka, Mueller (2014)
Tamborra, Hanke, Mueller, Janka, Raffelt (2013)

The next galactic SN may reveal distinct signatures of the neutrino mechanism

Fast Conversions = Heating



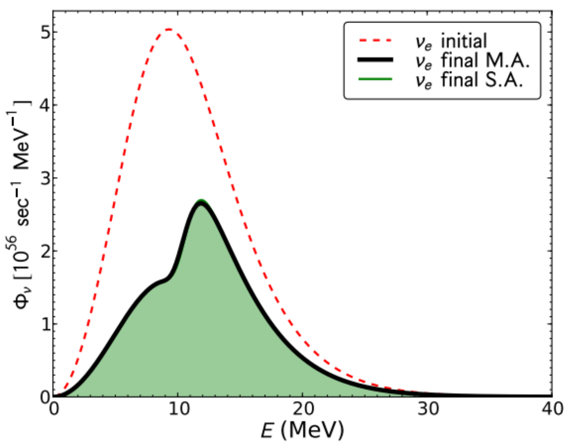


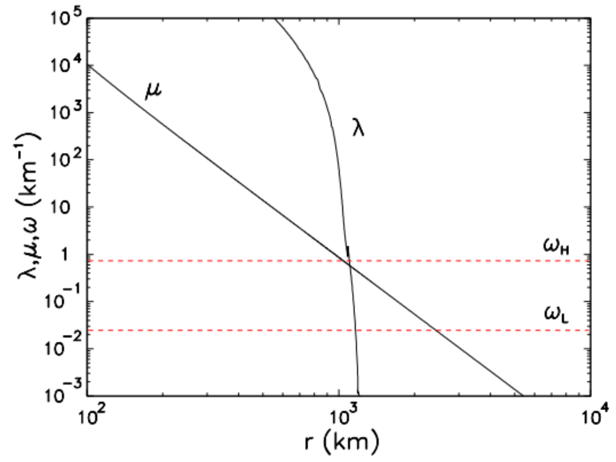
Not possible with slow conversions that occur above r_gain

May be possible if fast conversions that occur below r_gain!

Dasgupta, O'Connor, Ott (2011)

O-Mg-Ne Supernovae

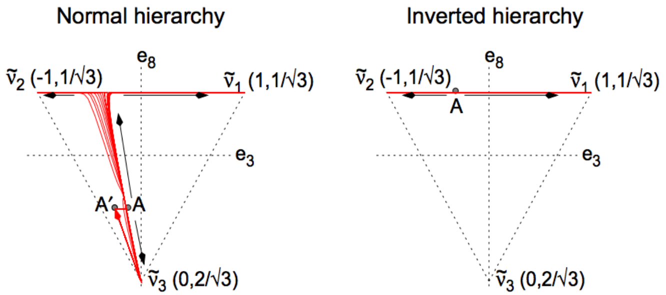




Duan, Fuller, Carlson, Qian (2007)
Duan, Fuller, Carlson, Qian (2008)
Dasgupta, Dighe, Mirizzi, Raffelt (2008)
Cherry, Fuller, Carlson, Duan, Qian (2010)

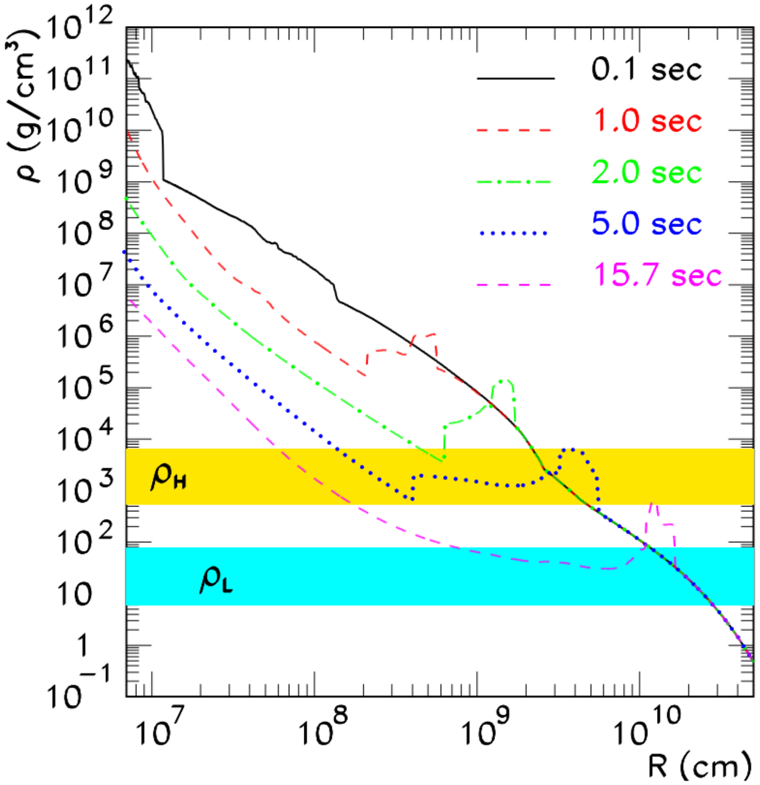
Star has sharply falling matter density, so MSW resonance occurs before (slow) collective effects.

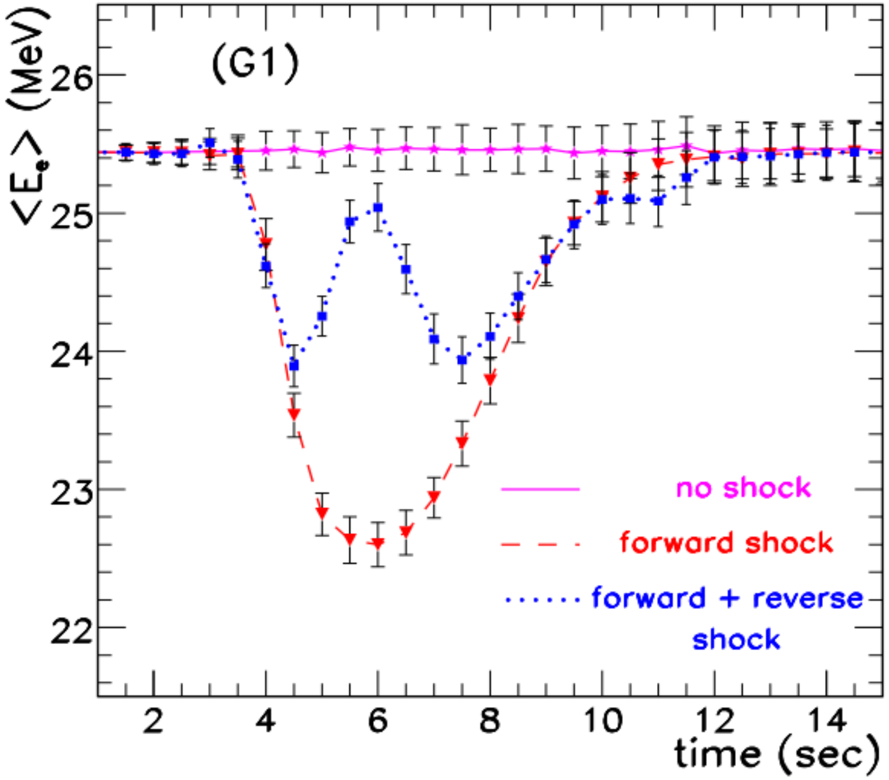
Explained using synchronized MSW



Application of 3 flavor formalism by Dasgupta and Dighe (2007)

Shock Wave

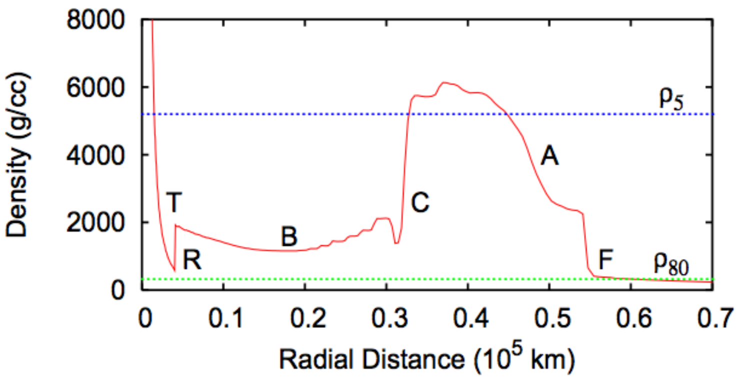


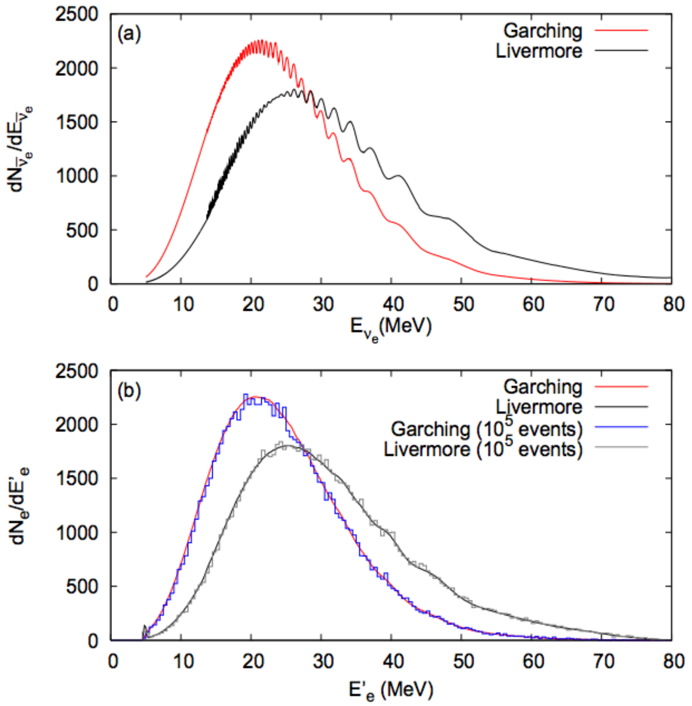


Schirato and Fuller (2002)
Tomas, Kachelriess, Raffelt, Dighe, Janka, Scheck (2004)

Shockwave propagation leads to dips/bumps in observed average E

Phase Effects



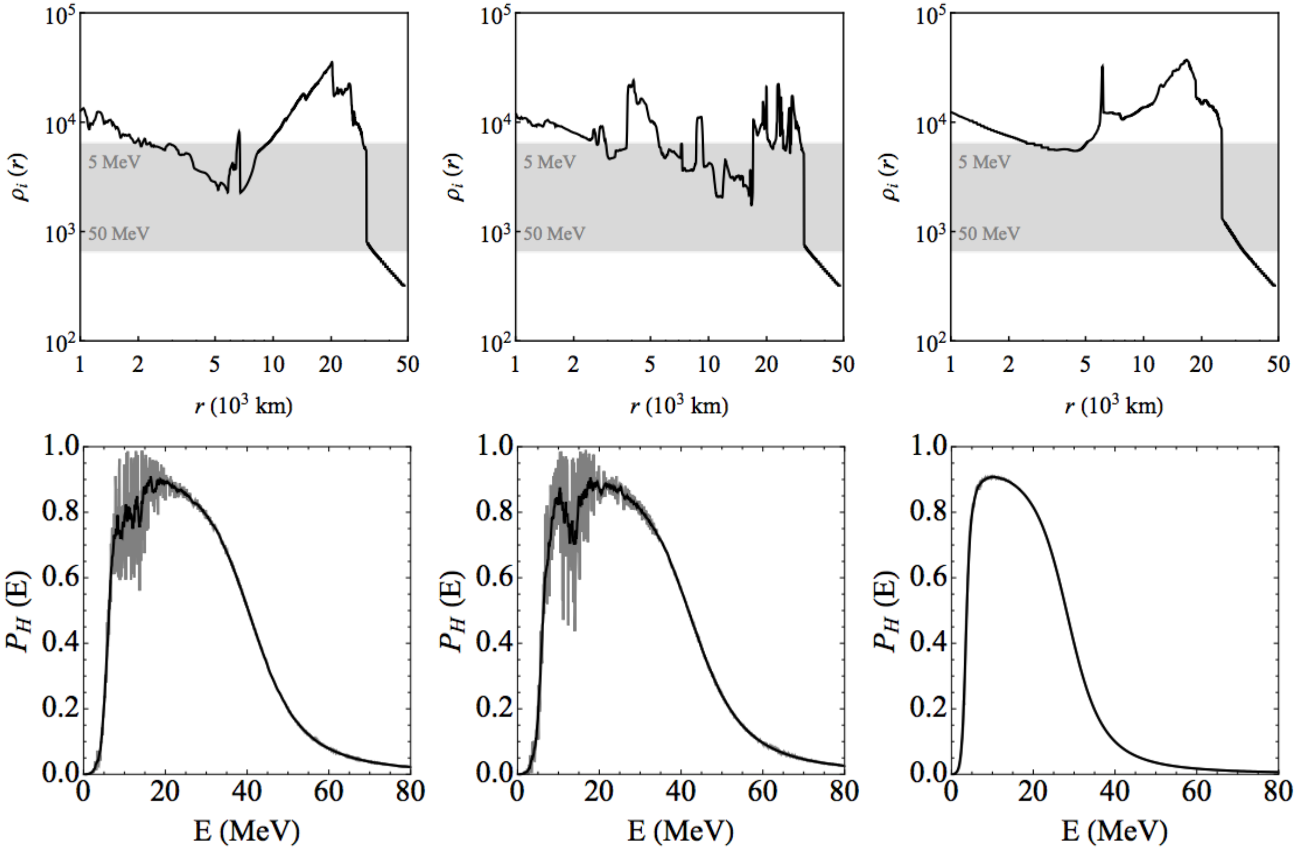


Density profile is not monotonic. There are multiple resonances. Some of them are close enough and lead to interference features with large wavelength in 1/E

Typically gets averaged out due to finite energy resolution

Dasgupta and Dighe (2005)

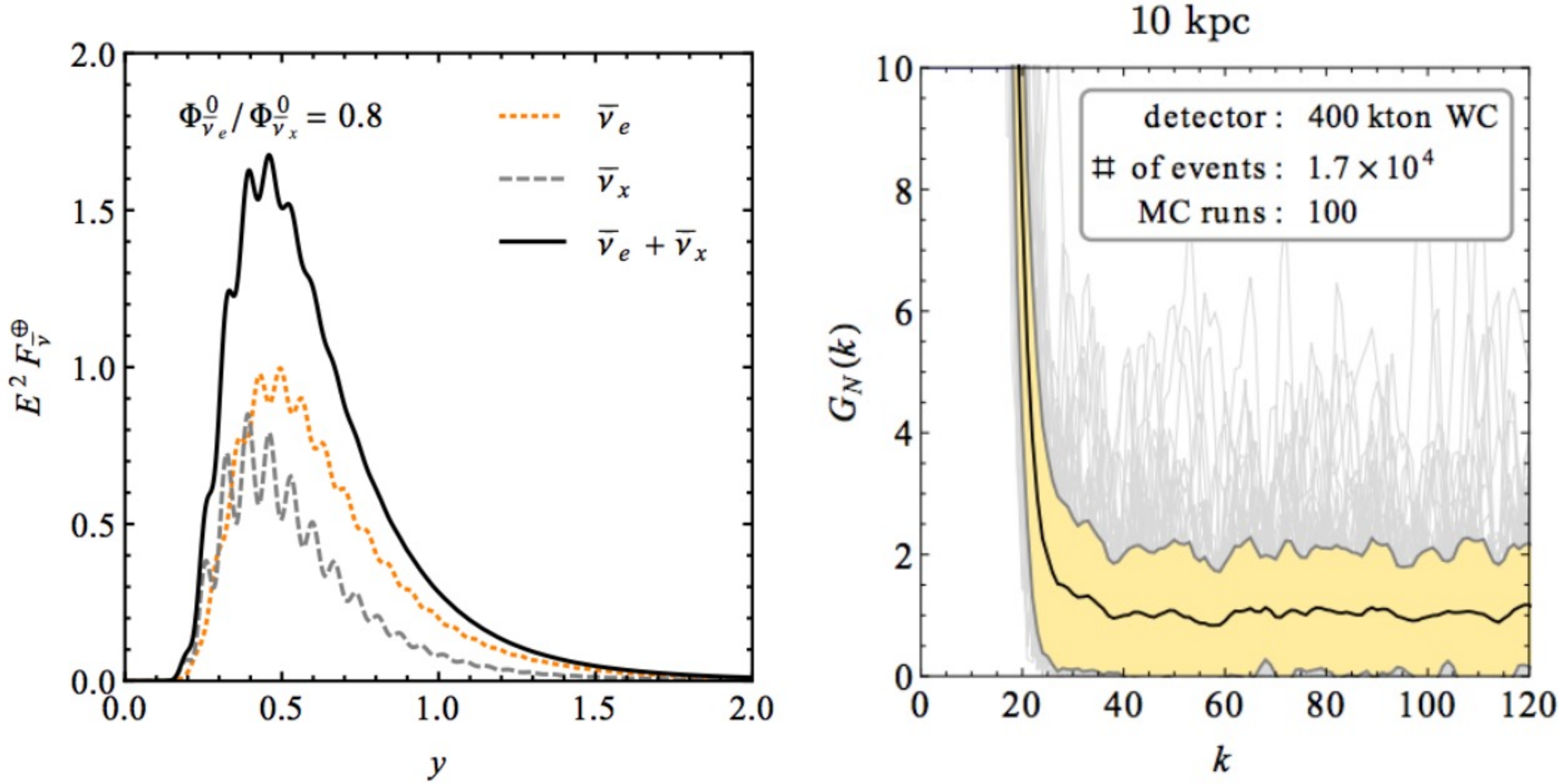
Turbulence



Plot from Borriello, Chakraborty, Janka, Lisi, Mirizzi (2013)
Fogli, Lisi, Mirizzi, Montanino (2006)
Friedland and Gruzinov (2006)

Survival probabilities are highly stochastic quantities during certain time-windows

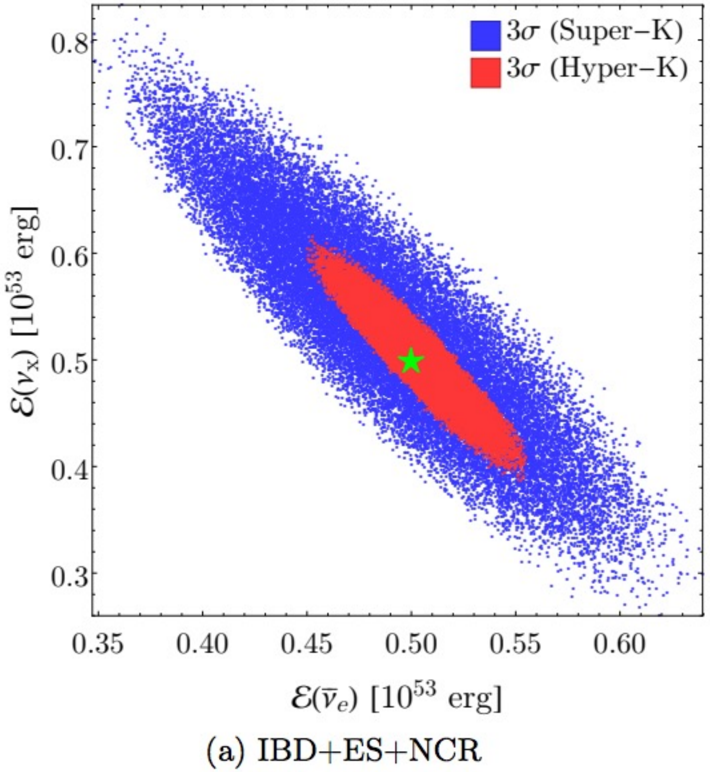
Earth Matter Effects



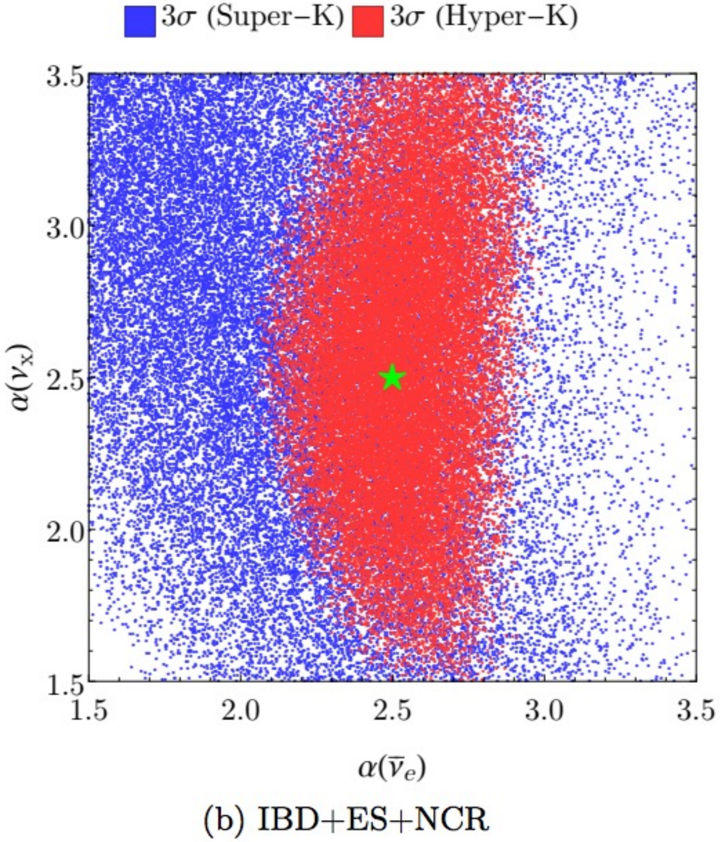
Energy-dependent regeneration in Earth depends on spectral differences between flavors and encodes neutrino mass ordering. May be hard to see.

Boriello, Chakraborty, Mirizzi, Serpico, Tamborra (2012)
Lunardini, Smirnov (2001)

SN Flux Reconstruction



(a) IBD+ES+NCR

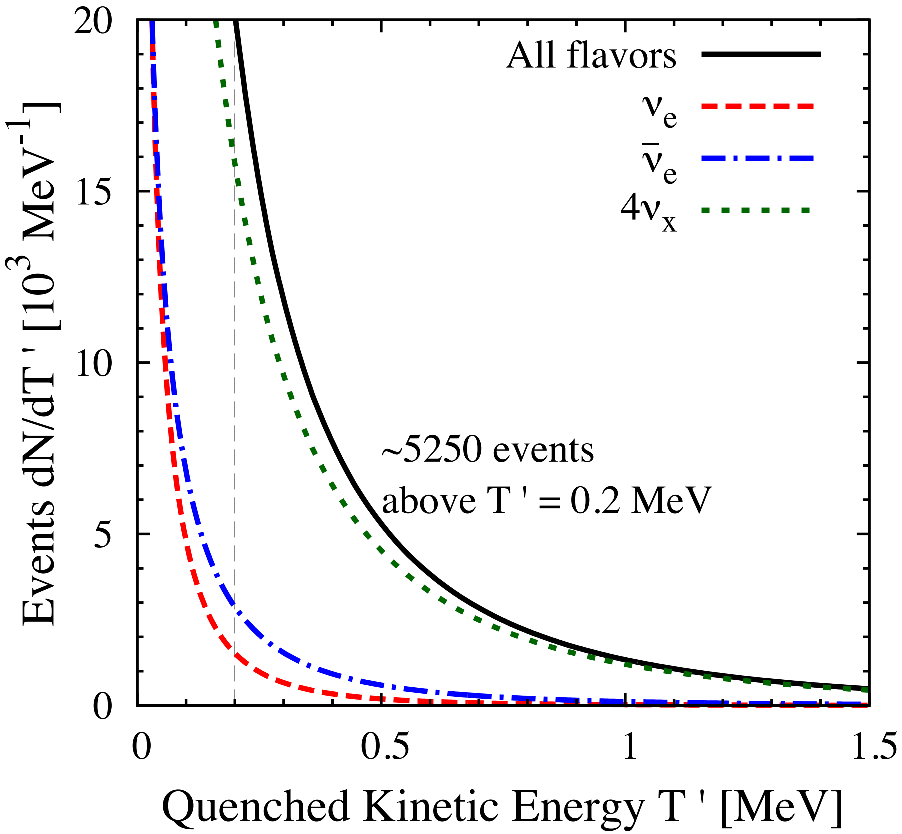


(b) IBD+ES+NCR

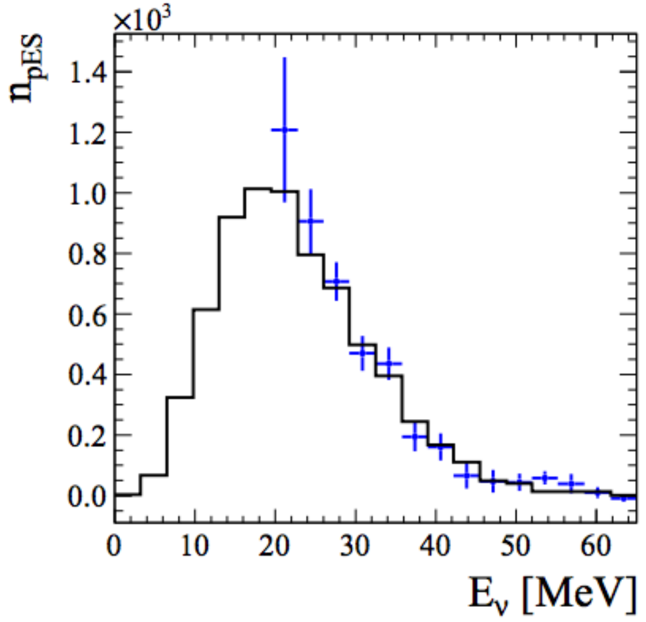
May be possible to reconstruct fluxes with unknown pinching

Rosso, Vissani, Volpe (2017)
For nue: Laha, Beacom (2014), Laha, Beacom, Agarwalla (2014), Nikrant, Laha, Horiuchi (2017)
Previously: Minakata, Nunokawa, Tomas, Valle (2008)

Neutral Current is Important



Neutrino – Proton elastic scattering can give the unoscillated fluxes if measured with enough statistics and reconstructed with precision



Beacom, Farr, Vogel (2003)
Dasgupta and Beacom (2011)
See also Chauhan, Dasgupta, Datar (2021) for a new idea using Deuterium

Detailed analysis for JUNO by Li, Li, Wang, Wen, Zhou (2017)

Some Reviews

Mirizzi, Tamborra, Janka, Saviano, Scholberg, Bollig, Huedepohl, Chakraborty
Riv. Nuovo Cimento 39 (2016)

Detailed SN Neutrino Review

Horiuchi and Kneller
J.Phys. G45 (2018)

Interpretative

Duan, Fuller, Qian
Ann.Rev.Nucl.Part.Sci. 60 (2010)

Chakraborty, Hansen, Izaguirre, Raffelt
Nucl.Phys. B908 (2016)

Shalgar and Tamborra - Annual Rev. (2021)
Capozzi, Saviano - Universe 8 (2022)
Richers and Sen - (2022)

Focussed on collective effects

Dasgupta
PoS ICHEP2010 (2010), and PoS NOW 2022 (2023)

Short reviews for the impatient

Takeaway

www.bdasgupta.com
bdasgupta@theory.tifr.res.in

- SN theory
  - Different phases of SN explosion and different oscillation physics
- Oscillation theory
  - Adiabatic and Non-adiabatic (at Shock) MSW
  - Slow/Fast Collective Mixing due to “Crossings”
- Detection Prospects, Bounds
  - Neutrinos of all Flavors, Energies, Times, Directions …
  - Multi-messenger helps
- Upshot
  - Neutronization burst : Physics is MSW-like
  - Accretion : If “Crossing” exists then Collective -> Mixing, else MSW-like
  - Cooling : Spectra are very similar; New and Improved BSM Bounds